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SKIN-FRICTION MEASUREMENTS IN INCOMPRESSIBLE FLOW

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SUMMARY

Experiments have been conducted to measure in incompressible flow the local surface-shear stress and the average skin-friction coefficient for a turbulent boundary layer on a smooth, flat plate having zero pressure gradient. The local surface-shear stress was measured by a floating-element skinfriction balance and also by a calibrated total head tube located on the surface of the test wall. The average skin-friction coefficient was obtained from boundary-layer velocity profiles. The boundarylayer profiles were also used to determine the location of the virtual origin of the turbulent boundary layer. Data were obtained for a range of Reynolds numbers from 1 million to about 45 million with an attendant change in Mach number from 0.11 to 0.32.

The measured local skin-friction coefficients obtained with the floating-element balance agree well with those of Schultz-Grunow and Kempf for Reynolds numbers up to 45 million. The measured average skin-friction coefficients agree with those given by the Schoenherr curve in the ranges of Reynolds numbers from 1 to 3 million and 30 to 45 million. In the range of Reynolds numbers from 3 to 30 million the measured values are less than those predicted by the Schoenherr curve.

The results show that the "universal skin-friction constants" proposed by Coles approach asymptotically a constant value at Reynolds numbers exceeding 21 million. Because of the scatter in the aforementioned constants and the limited Reynolds number range of the present investigation, there is some doubt as to the validity of any turbulent skin-friction law written on the basis of the present results. Hence, no new friction law is proposed.

The frictional resistance of a flat plate was calculated by means of the momentum method and

also the integrated measured local surface shear. For Reynolds numbers from 14 million to 45 million both methods give about the same result; whereas at lower values of Reynolds number the momentum method based on relocity profiles uncorrected for the effects of turbulence results in a frictional resistance as much as 4 percent higher than that of the integrated shear.

The measurement of local surface shear by a calibrated Preston tube appears to be accurate and inexpensive. The calibration as given by Preston must be modified slightly, however, to yield the results obtained from the floating-element skin-friction balance.

INTRODUCTION

In recent years there has been a resurgence of interest in the problem of the turbulent boundary layer on a smooth flat plate having zero pressure gradient along its length or breadth. This interest falls into two categories. First, it is necessary for the aeronautical designer to know the effect of Reynolds number variation on the average skin-friction coefficient for the accurate prediction of both drag and heat transfer. Secondly, there has been considerable controversy in England among hydrodynamicists with regard to the variation of average skin-friction coefficient with changing Reynolds number (see refs. 1 through 4) and hence the ability to project shipmodel test results to full-scale Reynolds numbers.

Up to this time much work has been done in attempts to determine a so-called skin-friction law for incompressible fluids. Most of this work has been experimental in nature, leading to a law having empirically determined constants.

It was intended in the present work to determine accurately the empirical constants required to write a skin-friction law by making use of the modern techniques now available for measuring local surface-shear stress and by use of extremely accurate manometers for measuring the local velocity in the boundary layer. By the use of such techniques, it was hoped that a friction law could be determined with an accuracy of ± 2 percent particularly for large Reynolds numbers. The investigation also included an evaluation of the accuracy of a method, proposed by Preston in reference 5, involving the use of a single surface tube to determine local surface-shear stress.

NOTATION

	NOTATION
C_f	local skin-friction coefficient, $\frac{ au_w}{q}$
C_F	average skin-friction coefficient, $\frac{2\theta}{x}$
C_p	pressure coefficient, $\frac{p-p_{\infty}}{q_{\infty}}$, dimen sionless
C_1	constant in skin-friction equation, $\frac{V\delta^*}{u^*\delta}$
C_2	constant in skin-friction equation, $ \frac{(\frac{V}{u^*})^2}{\delta} \frac{\delta^* - \theta}{\delta} $
d	inside diameter, in.
D	outside diameter, in.
II	shape parameter, $\frac{\delta^*}{\theta}$
k·	slope of wall law and velocity-defect law curves in the similarity region
M	Mach number
p	local static pressure, lb/sq in.
p_{w}	free-stream static pressure, lb/sq in.
p_t	local total pressure, lb/sq in.
$p_{t_{\infty}}$	free-stream total pressure, lb/sq in.
q_{ω}	free-stream dynamic pressure, lb/sq in.
R_x	Reynolds number, $\frac{V_x}{\nu}$
$R_{ heta}$	Reynolds number, $\frac{V\theta}{\nu}$
T	temperature, °F
u	local velocity, ft/sec
u^*	friction velocity, $\sqrt{\frac{\tau_w}{\rho}}$, ft/sec
1.	free-stream velocity, ft/sec
$w_{ m trip}$	weight flow of air ejected from boundary-layer trip, lb/sec
x	distance in the direction of flow from
m *	the virtual origin of the turbulent boundary layer, in.

vertical distance from wall, in. yspanwise distance across channel, ĩ measured from center line of chanboundary-layer thickness, δ at $\frac{u}{V} = 0.990$, in. boundary-layer displacement thick- δ^* ness, $\delta \int_0^1 \left(1 - \frac{\rho u}{\rho_\infty V}\right) d\left(\frac{y}{\delta}\right)$, in. boundary-layer momentum thickness, θ $\delta \int_{0}^{1} \frac{\rho u}{\rho_{\infty} V} \left(1 - \frac{u}{V}\right) d\left(\frac{y}{\delta}\right)$, in. absolute viscosity, lb sec/sq ft kinematic viscosity, sq ft/sec local density, lb sec²/ft⁴ free-stream density, lb sec²/ft⁴ surface friction stress, lb/sq ft $\frac{u}{u^*}$ when $\frac{yu^*}{v}$ =1.0 or when $\log_{10} \frac{yu^*}{v}$ = $\phi(0)$ $\phi(1) - \phi(0) = \frac{0 \text{ (See fig. 15(a).)}}{V - u} \text{ when } \frac{yu^*}{\delta^* V} = \frac{1}{C_1} \text{ (See fig. 15(b).)}$ **EQUIPMENT**

MODEL

The friction measurements were made on a flat plate which formed one wall of a channel mounted in the wind tunnel as shown in figure 1. The test wall was mounted between a pair of end plates to which was attached an adjustable auxiliar, plate approximately parallel to the test wall. The auxiliary plate could be adjusted to change the longitudinal pressure gradient in the

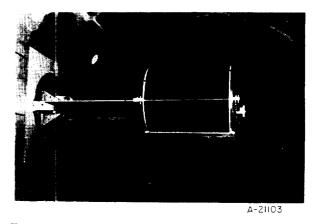


FIGURE 1.—Boundary-layer channel in the test section of the Ames 12-foot pressure wind tunnel.

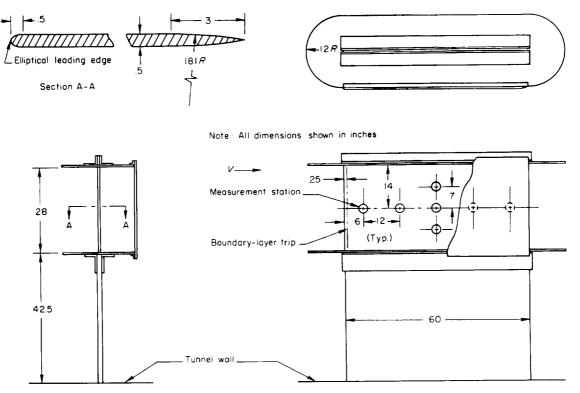
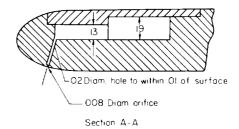


FIGURE 2.—-Three-view drawing of boundary-layer channel.

channel. Preliminary measurements indicated that without the auxiliary plate the longitudinal pressure gradient along the test wall was not uniform.

The test wall of the channel and the auxiliary wall opposite were identical in cross section. The nose was elliptical with a ratio of major axis to minor axis of 2.0. The trailing edge was sharp, having a circular-arc section tangent to the surface 3 inches forward of the trailing edge (fig. 2). The test wall was made of mild steel polished to a fine finish. Measurements with an interferometer indicated that, generally, the test wall had a surface finish of 20 to 40 microinches (peak to valley). There were a few streamwise scratches on the surface which were deeper than this but it is believed that they had little or no effect on the flow.

The other three walls of the channel were made of aluminum and had a finish about equal to that of the test wall. All holes and joints were sealed to prevent the flow of air from the higher pressure stream of the tunnel into the channel at other than the front opening. A permanent boundary-layer trip was installed near the leading edge of the test wall (fig. 3). This trip was of the air ejection type used by



Note: All dimensions shown in inches

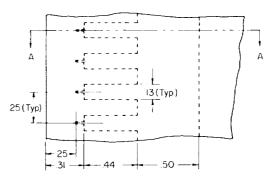


FIGURE 3.-- Details of boundary-layer trip.

Fage and Sargent (ref. 6). The trip will be discussed further in the section on test conditions.

WIND TUNNEL

This experimental investigation was done in the Ames 12-foot pressure wind tunnel. The wind tunnel is of the variable-density type providing Reynolds numbers up to 10 million per foot at a Mach number of about 0.30 and Reynolds numbers up to 1.7 million per foot at Mach numbers up to about 0.95. The turbulence level of the wind-tunnel air stream is very low.

EXPERIMENTAL METHODS

The reliability of skin-friction measurements is critically dependent upon the precision of the measuring apparatus. It, therefore, seems appropriate to discuss in some detail the characteristics of the measuring apparatus, the degree of precision attained, and the procedure used in conducting the tests.

LOCAL VELOCITY MEASUREMENTS

The velocity profiles through the boundary layer were determined from measurements with a total-pressure tube and a static-pressure orifice in the plate, located at the same longitudinal station. The total-pressure tube was very carefully constructed with a flattened end which was 0.007-inch high and 0.080-inch wide. The wall thickness was 0.002 inch. (See fig. 4.) The opening of the tube was perpendicular to the direction of the free stream and was free of burrs and imperfections. This tube was mounted on a screw device which allowed it to be moved perpendicular to the wall. This screw was calibrated and found to be capable

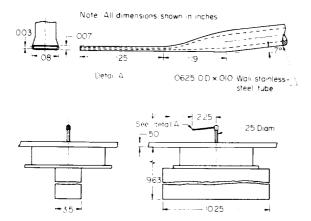


Figure 4.—Details of velocity probe and probe mounting mechanism.

of positioning the tube to 0.001 inch. The zero position of the tube was determined by an electrical circuit which was energized when the total-head tube made contact with the plate. This method was quite successful and was found to be capable of consistently indicating the zero position to 0.001 inch. This accuracy was only possible if the wall and tube were kept scrupulously clean and free of all oxides, oil, and foreign matter.

The quantities measured were the local total pressure in the boundary layer, the static pressure at the wall, and the vertical distance from the surface of the wall to the center line of the face of the total-pressure tube. It was assumed that the static pressure was constant through the boundary layer and that the total temperature in the boundary layer was equal to the total temperature in the tunnel settling chamber. Because of the small vertical dimension of the total-pressure tube, no correction was applied to the measured height of the tube above the test wall to account for the apparer t displacement of the tube resulting from the total-pressure gradient through the boundary layer. No correction was applied to the velocity profiles for the effect of turbulence.

An a Iditional probe was constructed for use in the determination of the location of boundarylayer transition. The longitudinal variation of the surface velocity near the leading edge of the plate was measured. This device was capable of traversing the plate in a streamwise direction from 0.5 inch aft of the leading edge to about 3.25 inches aft of the leading edge. The local total pressure was measured with a probe having the same dimensions as the one previously described (see fig. 4) and the static pressure was measured with a 0.035-inch-diameter static-pressure probe located 1.0 inch away from the surface of the plate and at the same longitudinal station as the total-pressure probe. Local velocities were computed from these measurements using the same assumptions as were made for the surveys through the boundary layer.

LOCAL SURFACE-SHEAR MEASUREMENT

The local surface-shear stress was measured by two different techniques. The first of these made use of a floating-element device which measured the shear stress directly. The second technique made use of a calibrated total-pressure tube mounted on the surface of the wall as proposed by J. H. Preston. Preston made measurements with

air flow in a pipe, whereas the present measurements with surface tubes were made to validate and determine the accuracy of the technique for air flow on a flat plate and to verify Preston's calibration of the tubes.

Floating-element device.—The local surfaceshear stress was measured by a floating-elementtype device similar to that used by Dhawan (ref. 7) and others. The floating-element technique was also used by Schultz-Grunow (ref. 8) and Kempf (ref. 9) in their historically important surfaceshear measurements.

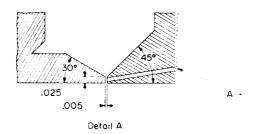
Since little is known about the effect of change of the size of gap around the floating element on the measured surface shear, it was decided to construct a device whose element could be repositioned and centered in the gap. Both Schultz-Grunow and Kempf used such a device while Dhawan and others used a simple deflection-type instrument. In the present unit the floating element was repositioned by a small, powerful electromagnet. The position of the element was indicated by a differential transformer capable of indicating movement of the floating element to an accuracy of a few millionths of an inch. When the position indicator showed that the floating element had started to move from its no-load neutral position, the strength of the electromagnet was varied until the element returned to its no-load neutral position. Since the electromagnetic force was equal and opposite to the drag force exerted on the element, the average surface-shear stress on the floating element could be deduced from the measured electromagnetic force and a predetermined calibration.

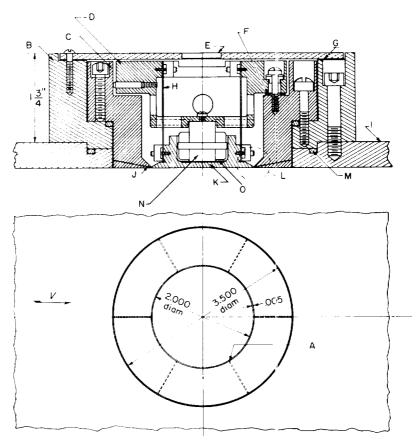
The shear-stress measuring device was capable of indicating the drag force on the element with a sensitivity of about 0.02×10^{-3} pounds for a range of force from 0 to about 30×10^{-3} pounds. The accuracy of determining the load under test conditions is believed to be within ± 2 percent of applied load throughout the load range encountered in the tests. Calibrations of the element displayed extremely good repeatability. The measured data were corrected for effects of change in temperature of the unit.

In figure 5 is presented a detailed drawing of the shear-stress measuring device. This device consisted of a 2-inch-diameter plate which was mounted on very limber flexure pivots. The flexure pivots were, in turn, attached to a sturdy support frame which was mounted on the working wall of the boundary-layer channel. As may be seen in figure 5, the support frame and movable plate were mounted on the channel wall in an integral unit. The 2-inch-diameter movable plate was centered in a 2.010-inch-diameter hole in the support frame with its working surface set flush with the working surface of the support frame. The surface of the floating-element unit was carefully alined flush with the surface of the channel wall using both dial and interferometric indicators. It was possible to position the element surface within about ± 0.00005 inch by means of the dial indicator.

Tests were made to study the effects of small variations in flushness of the floating element with the surrounding fixed surface. Measurements of surface shear at identical test conditions were made for a range of positions of the floating element, both depressed below and protruding above the fixed surface of the plate. It was found that the surface of the floating element could be depressed as much as 0.0005 inch without any change in the surface shear. However, when the element protruded above the surface of the wall, there were noticeable deviations in the measured shear force. Consequently, the surface element was always maintained flush with or slightly below the surface of the channel wall.

The entire floating-element unit was constructed of Invar in order to minimize the effect of temperature changes on the calibration of the unit. The faces of both the support frame and the floating plate were very carefully lapped to ensure both a fine surface finish and also flat surfaces having sharp edges on the inside and outside diameters of Interferometric measurements indithe units. cated that the surface had a peak-to-valley roughness of about 10 to 20 microinches and a flatness of about 20 to 40 microinches. The floating-element unit was made into a pressure-tight capsule to prevent the flow of high-pressure air from the tunnel main flow into the higher speed flow of the boundary-layer channel. Damping of the floating element was achieved by using 20,000 centistoke oil in a cup machined integral with the back of the element. The cup was adjusted to have 0.005inch clearance with the displacement indicator and electromagnet which are fixed to the support frame (see fig. 5).





- A -- Pressure orifice
- B -Mounting ring
- C -Support frame
- D—Support for displacement indicator and electromagnet
- E—Electrical connector mounting hole
- F—Spring washer
- G -Cover pressure seal
- H-Flexure

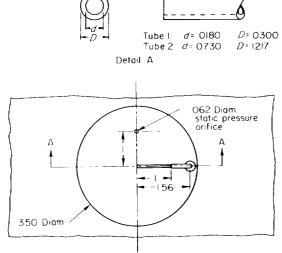
- t -Chani el wall
- J -Tube from gap pressure measuring orifice
- K -Floating element
- L Detai A
- M —"O" ring seal
- N -Displacement indicator and electromagnet
- O—Damping fluid—Dow Corning 200 fluid 20,000 centistokes

FIGURE 5. Details of the floating-element skin-friction balance.

The static pressure in the gap between the floating element and the channel wall was measured by means of six static orifices in the gap (see fig. 5) and a buoyancy correction was applied to the surface-shear force measured on the element. This correction was always less than 1 percent of the applied force on the element.

Surface-tube shear-stress device.—In 1953 a very simple technique for measuring surface-shear stress was proposed by J. H. Preston (see ref. 5). This technique made use of the total pressure measured by a round total-head tube mounted flush with the surface (see fig. 6). The pressure measured by the total-head tube in conjunction with the surface static pressure measured at the same location along the plate was calibrated by Preston in terms of the local surface-shear stress.

Two total-head tubes having outside diameters of 0.0300 inch and 0.1217 inch, respectively, were



Note: All dimensions are shown in inches

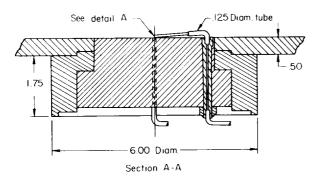


Figure 6.—Details of the Preston surface-shear tubes. 509374—59——2

used in the present investigation. The tubes had a ratio of inside diameter to outside diameter of 0.600, the same proportions used by Preston. Care was taken to make the mouth of the tube perpendicular to the longitudinal axis of the tube. The equations and assumptions used in the reduction of the measured data are given in reference 5.

SENSITIVE MANOMETER

In order to measure the velocities in the boundary layer and the pressures associated with the surface-tube shear-stress device with sufficient accuracy to give an over-all accuracy of results of 1 percent or better, it was necessary to devise a manometer capable of measuring very small pressure differences over a large range of pressure difference. Such a device was designed and built and was found to be capable of indicating a pressure difference of about 0.06 pound per square foot with an accuracy of ± 0.12 pound per square foot for pressure differences as large as 600 pounds per square foot.

This manometer was of the U-tube type with a float in the low-pressure leg of the system. This float had a steel slug incorporated in it and a servo-operated follower mounted on a lead screw alongside the manometer leg to indicate the position of the slug in the leg of the manometer. The lead screw was calibrated in terms of the pressure difference applied across the two legs of the manometer. The glass tubes used in this manometer were precision bored to have an inside diameter of 0.750 ± 0.001 inch. The fluid (tetrabromoethane) in the manometer was maintained at a fixed temperature of 107° F $\pm \frac{1}{2}^{\circ}$ F.

The bore of each manometer tube was coated with Dri Film, a General Electric silicone product, to reduce the effect of the meniscus of the fluid on the pressure readings. Calculations show that the capillary effect of the meniscus could result in a maximum error of about 0.2 pound per square foot in the pressure reading if the angle of contact between the manometer fluid and the glass tube varied from 0° to 90°. Because of the Dri Film coating it is felt that the error in measured pressure due to capillary forces has been reduced to a value considerably smaller than the accuracy of the indicating system of the manometer.

Due to the fact that this instrument had a large range of indication and extremely high sensitivity, the calibration of the instrument posed some difficulty. Since there was no instrument available to use as a standard, it was decided to determine the specific gravity of the fluid at the stabilized temperature (107° F) and use this as the calibration of the instrument in conjunction with an accurate calibration of the lead screw follower.

TEST CONDITIONS

The Reynolds number in the present tests varied from about 1 million to 10 million per foot of channel length. This range of Reynolds numbers was obtained by varying the tunnel total pressure from 8 to 80 pounds per square inchabsolute and the Mach number from 0.11 to 0.32. These values of Mach number are in the range where compressibility effects in the air flow are generally considered insignificant.

VELOCITY PROFILES

The boundary-layer velocity profiles were measured at stations 1.312, 2.312, 3.312, and 4.312 feet aft of the leading edge of the channel wall. The most forward measurement station (0.312 feet aft of leading edge) was not used since the velocity profiles were distorted and were of no interest. The longitudinal locations used provided Reynolds numbers based on the distance

from the leading edge from about 1 million to about 43 million.

SURFACE-SHEAR STRESS

Local surface-shear stress was measured at stations 1.5, 2.5, 3.5, and 4.5 feet aft of the leading edge of the channel wall as is shown in figure 2. Again the most forward measuring station (0.5 feet aft of the leading edge) was not used because of the distorted velocity profiles. The Reynolds number based on the distance of these stations from the leading edge varied from about 1.5 million to 45 million.

LONGITUDINAL PRESSURE GRADIENT

The longitudinal static-pressure gradient measured on the test surface of the boundary-layer channel is presented in figure 7. At the leading edge of the channel there was a pressure peak which is not shown in the figure. Throughout the major portion of the channel, where measurements were made, the local static pressure varied less than about 0.5 percent of the velocity head from the reference static pressure at the longitudinal midpoint of the channel. As may be seen in figure 7 there was little effect of change in either Mach number or tunnel total pressure on the pressure gradient.

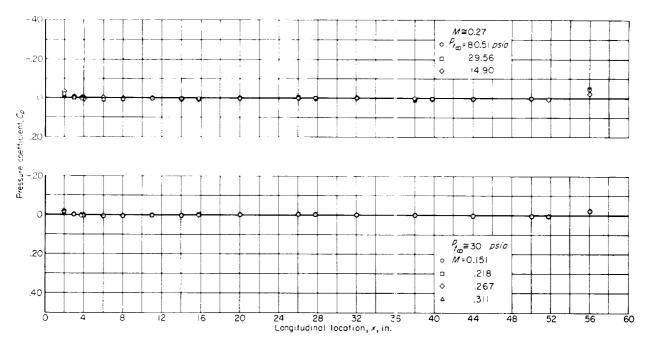


FIGURE 7.—Longitudinal pressure gradient in boundary-layer channel.

BOUNDARY-LAYER TRIP

A boundary-layer trip was provided to assure a two-dimensional turbulent boundary layer near the leading edge of the working wall of the channel. An air-injection-type trip was chosen because it could be readily varied in strength to trip the boundary layer with the least amount of disturbance. The geometry of the trip is given in figure 3.

The quantity of air to be ejected from the trip was determined using longitudinal velocity surveys which were made at the surface of the plate with a total-pressure and a static-pressure tube. Typical longitudinal velocity distributions at the plate surface for various amounts of ejected air are presented in figure 8. When no air was ejected

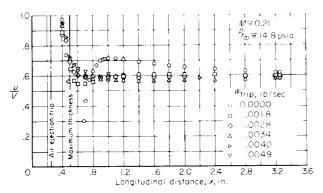


FIGURE 8.—Effect of varying quantity of boundary-layer-trip air on local velocity near channel wall.

from the trip, it appeared that some type of separation phenomenon was present. However, when air was ejected from the trip, this phenomenon disappeared and it seemed that the boundary layer became turbulent within about 0.25 inch of the trip. It was not possible to keep the probe on the surface of the wall forward of the maximumthickness point and therefore the data forward of this point do not represent surface measurements.

For the Mach number and total-pressure condition in figure 8, the air quantity selected as that which assured a turbulent boundary layer with the least distortion was 0.0034 pound per second. A similar set of surveys was made for each test condition and the air quantities selected in this manner were utilized for their particular test conditions.

TWO-DIMENSIONALITY OF FLOW

As was previously mentioned the walls of the channel were capable of being moved with respect to one another to provide for adjustment of the longitudinal static-pressure gradient. These walls were also adjusted so that the static pressure did not vary in the transverse direction.

To check the two-dimensionality of the flow, boundary-layer velocity profiles were measured at three spanwise locations at the same longitudinal station. The spanwise locations chosen for the measurements were at the center line of the working wall of the channel and at 7 inches either side of the center line. These three profiles for several test conditions are presented in figure 9. Their

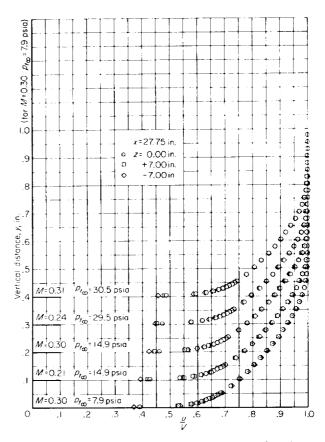


Figure 9.—Effect of spanwise location in boundary-layer channel on velocity profile.

similarity indicates a flow which closely approximates two-dimensional flow.

DETERMINATION OF VIRTUAL ORIGIN OF TURBULENCE

Physically, the turbulent boundary layer can not start with zero thickness and the virtual origin of the turbulent layer must therefore be estimated. One simple method for making such an estimate was proposed by Rubesin, et al. (ref. 10), and this method has been used in the present report.

The virtual origin of turbulence was estimated by plotting $\log 2\theta$ versus $\log x$ (where x is the distance from the leading edge of the test surface) and determining the magnitude of the change in x required to make the slope of the line equal to some reference value.

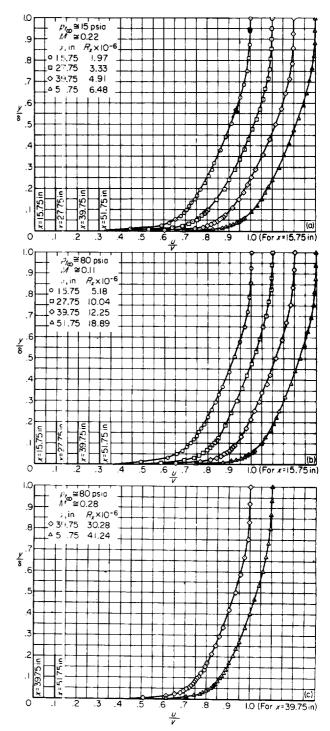
The reference value of the slope, d (log 2θ)/d (log x), which was used was the mean value of the slopes computed for each of four logarithmic laws presented in reference 11. (The law by Schultz-Grunow was omitted.) The reference value of slope used for the estimation of the virtual origin varied from about 0.826 to 0.850 for a variation of Reynolds number per foot from about 1 million to 10 million.

It was found that for all conditions at which tests were made, the change in x was within ± 1 inch and in many cases was within $\pm \frac{1}{2}$ inch. On the basis of this analysis and due to the fact that the results scattered on both sides of zero, it was concluded that the leading edge of the working wall of the channel could be used as the virtual origin of the turbulent boundary layer and the distance from the leading edge to the point of measurement could be used as the reference distance for Reynolds number.

PRESENTATION AND DISCUSSION OF RESULTS

The principal results of the investigation are presented in tables I, II, and III. Table I contains measured velocity profile data for all test conditions. Table II contains the measured values of local skin-friction coefficient as a function of Reynolds number. In table III is presented a summary of the major boundary-layer parameters obtained from the boundary-layer velocity profiles.

There are presented in figure 10 some of the velocity profiles tabulated in table I. These profiles are typical of the profiles obtained for all test conditions. All of the measured velocity profiles have been mechanically integrated to obtain both the boundary-layer displacement thickness, δ^* , and momentum thickness, θ . The ratio of these two parameters, known as the shape parameter, II, has been computed and tabulated in table III. These results are presented in figure 11 as a function of the Reynolds number. As was expected



- (a) $p_{t\infty} = 15$ psia, M = 0.22
- (b) $p_{t\infty} = 80$ psia, M = 0.11
- (c) $p_{t\infty} = 80$ psia, M = 0.28

FIGURE 10. -Boundary-layer velocity profiles.

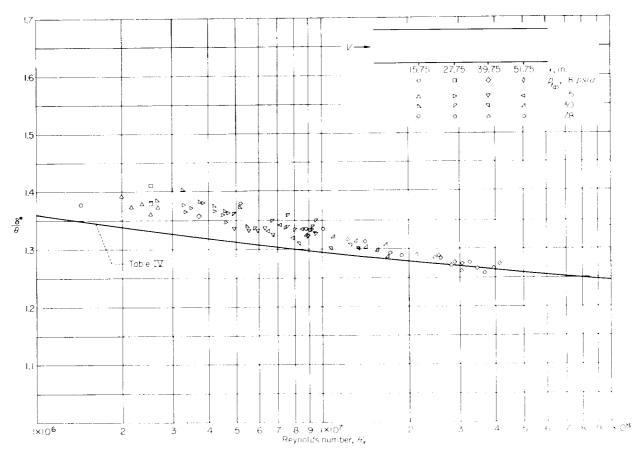


Figure 11. Variation of shape parameter with change in Reynolds number, R_z .

the shape parameter decreased as the Reynolds number increased. There appears to be large scatter in the data but this is not surprising since it is very difficult to obtain accurate values for either δ^* or θ . The line identified as table IV in this figure and in figures 12 and 13 will be discussed in a subsequent section.

The variation of the average skin-friction coefficient with change in Reynolds number is presented in figure 12. The average skin friction, C_F , was computed using the momentum thickness obtained from integration of the velocity profiles (presented in table III) measured at several stations along the wall of the channel. The Schoenherr line obtained from reference 11 will be discussed in a subsequent section.

There are presented in figure 13 the results of the measurement of the surface-shear stress. These results are tabulated in table II. The surface-shear stress was measured by the floatingelement technique previously described.

COMPUTATION OF DRAG BY MOMENTUM DEFECT AND BY INTEGRATION OF LOCAL SKIN FRICTION

The friction drag of a surface can be computed by two methods. The first of these methods involves computation, by mechanical integration of the boundary-layer profile, of the loss of momentum in the boundary layer which is directly convertible to the drag loss (data of fig. 12). The second method consists of integration of the local surface shear along the surface which is also directly convertible to the drag loss (data of fig. 13). A difficulty is involved in the second method in that it is necessary to know the local skin friction right up to the origin of the turbulent boundary layer. To circumvent this problem in the present investigation the drag at a point 18 inches aft of the leading edge of the surface of the channel was assumed to be that obtained by the momentum defect method. The local skin friction was then integrated and added to the assumed value of drag which resulted in a total drag at a particular

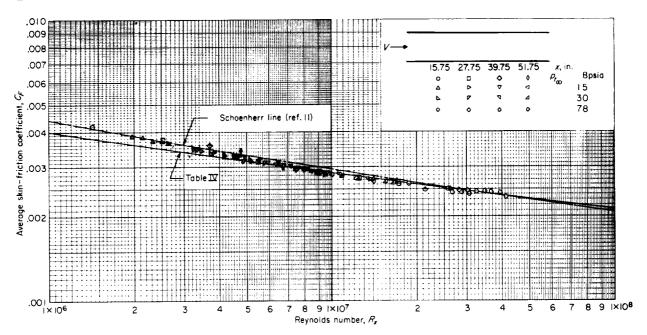


Figure 12. Variation of average skin-friction coefficient with change in Reynolds number, R_x ; boundary-layer velocity profile measurements.

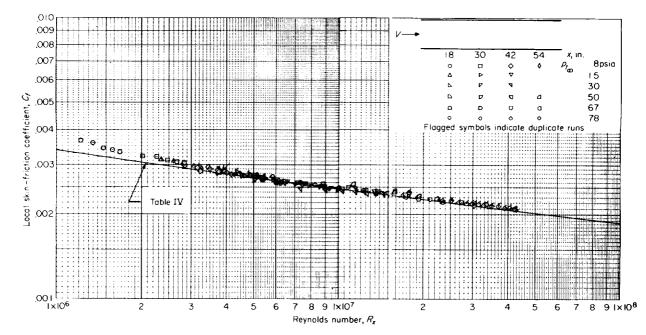
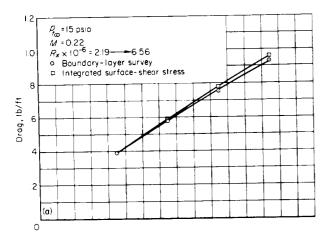
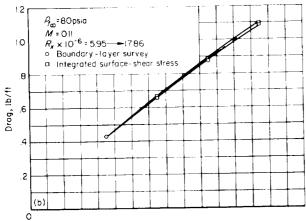


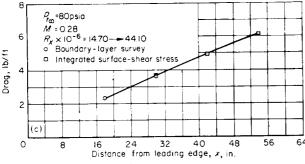
FIGURE 13. Variation of local skin-friction coefficient with change in Teynolds number, R_x ; floating-element technique,

longitudinal position on the channel wall. There are presented in figure 14 the results of these computations. The drag obtained by the momentum defect method is compared with that obtained by the integration of the local surface-shear stresses. Again it is pointed out that the drag at a point

18 inch is aft of the leading edge is assumed to be the same for both methods. It is apparent that at the smaller values of Reynolds number there is a discrepancy between the drag obtained by the two methods. At a Reynolds number of about 6.5 million the drag obtained from the integrated







- (a) $R_x \times 10^{-6} = 2.19$ to 6.56
- (b) $R_x \times 10^{-6} = 5.95$ to 17.86
- (e) $R_x \times 10^{-6} = 14.70$ to 44.10

Figure 14.—Comparison of drag computed by both the momentum defect and integrated surface-shear methods.

surface shear is about 4 percent higher than that obtained by the momentum defect, while at the highest Reynolds number of about 44 million the discrepancy between the two drags is reduced to almost zero.

METHOD OF DATA ANALYSIS

The aforementioned data will be discussed further in conjunction with a method of boundary-layer analysis previously used by Coles and others and described in some detail in reference 12. It is not felt that a detailed reiteration of the method is necessary here. The use of this method facilitates the analysis of the data of the present investigation in a systematic manner.

The equations of reference 12 which are used in the present analysis are given below in the notation of this report.

$$\begin{split} C_{f}R_{x} &= 2C_{1}e^{-k\varphi(1)}e^{k\sqrt{2/C_{f}}}\bigg[1 - \bigg(\frac{2}{k} + \frac{C_{2}}{C_{1}}\bigg)\sqrt{\frac{C_{f}}{2}} \\ &+ \frac{1}{k}\bigg(\frac{1}{k} + \frac{C_{2}}{C_{1}}\bigg)C_{f}\bigg] \\ C_{F}R_{x} &= 2C_{1}e^{-k\varphi(1)}e^{k\sqrt{2/C_{f}}}\bigg(1 - \frac{C_{2}}{C_{1}}\sqrt{\frac{C_{f}}{2}}\bigg) \\ \frac{\delta^{*}}{\theta} &= \frac{1}{1 - \frac{C_{2}}{C_{1}}\sqrt{\frac{C_{f}}{2}}} \\ R_{\theta} &= \frac{C_{F}R_{x}}{2} \end{split}$$

The analysis depends on the evaluation of the parameters k, $\varphi(1)$, C_1 , and C_2 which appear in the above equations.

The first step in the analysis is to express the velocity profiles in terms of the "law of the wall" $[u/u^* = f(yu^*/\nu)] \ \ \text{and the "velocity defect law"} \\ [(V-u)/u^* = f(yu^*/\delta^*V)]. \ \ \ \Lambda \ \ \text{typical profile in}$ terms of the "wall law" is presented in figure 15(a) while the same profile in terms of the "velocity defect law" is presented in figure 15(b). As may be noted on these figures both curves have a linear region when plotted on a semilogarithmic basis. From a comparison of the slopes of the linear portions of these curves it appears that they both have the same value. This portion of the curves is known in the literature as the region of overlap of the two laws or the region of similarity of the boundary layer. The existence of this region of similarity makes it possible to analyze the turbulent boundary layer quite readily. With the velocity profiles in this form the parameters k and $\varphi(1)$ may be evaluated. The parameter kis the slope of the curves in the similarity region. The parameter $\varphi(1)$ is the sum of the value of $\varphi(0)$ obtained from the wall law as shown in figure 15(a) and the value of $\varphi(1) - \varphi(0)$ obtained from the velocity defect law as shown in figure 15(b). The two parameters C_1 and C_2 are obtained from the velocity profile parameters as indicated by the definition given in the Notation section. The values of the four parameters k, $\varphi(1)$, C_1 , and C_2 may then be inserted in the skin-friction equations given previously to calculate a frictional resistance law for a fully developed turbulent boundary layer which starts at some point with zero thickness and grows as a fully developed turbulent boundary layer.

The variation of the parameters k, $\varphi(0)$ and $\varphi(1) - \varphi(0)$ with Reynolds number, R_{θ} (based on the boundary-layer momentum thickness), is shown in figure 16, while the variation with R_{θ} of C_1 and C_2 is shown in figure 17. It is expected that these turbulent boundary-layer parameters will become independent of Reynolds number if they are determined from measurements at large enough Reynolds numbers on an aerodynamically smooth plate in flow having zero pressure gradient. This appears to be the case in the present experiments for Reynolds numbers, R_{θ} , greater than about 26 thousand or a Reynolds number, R_x , of about 21 million. The average value of the constants in the range of Reynolds number independence were used in conjunction with the skin-friction equations given previously to make calculations of a frictional resistance law. The values of the constants used in this calculation were:

$$\varphi(1) - \varphi(0) = 3.00$$

$$\varphi(0) = 7.15$$

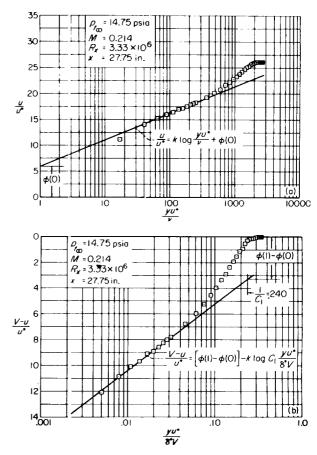
$$k = 5.00$$

$$C_1 = 4.00$$

$$C_2 = 25.9$$

The results of this calculation are presented in table IV.

As a result of the scatter in the values of these parameters, which were obtained from the experimental data, and the limited Reynolds number range attained in this investigation, there is some doubt as to the absolute values of the parameters listed above. Hence, a new frictional resistance law is not being proposed although the results of the calculation have been tabulated and presented in this form to afford a basis of com-



(a) Wall law.

(b) Velocity defect law.

Figure 15.—Boundary-layer velocity profiles in terms of the "vall law" and the "velocity defect law."

parison between the measured data of the present investigation and those of previous investigations.

COMPARISON OF COMPUTED FRICTION LAW WITH MEASURED DATA

The results of the frictional resistance law calculation presented in table IV are also presented in figures 11, 12, and 13. As was previously stated, constants applicable only in the range of Reynolds number above about 21 million were used in this calculation.

Shape parameter.—As may be seen in figure 11, the computed values of the shape parameter, δ^*/θ presented in table IV represent those computed from the measured velocity profiles only at the highest Reynolds numbers. This is not difficult to understand when it is realized that the computed value of the shape parameter

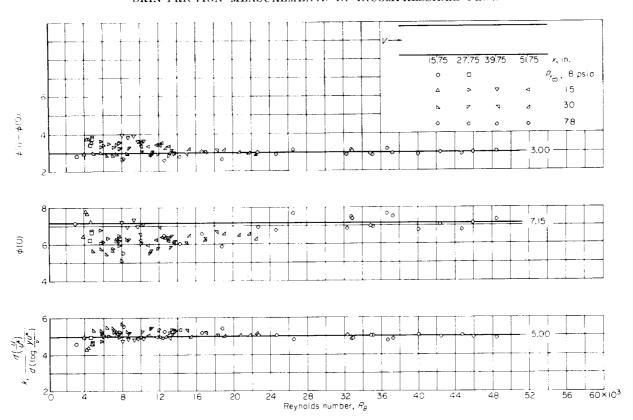


Figure 16. Variation of K, $\varphi(0)$, and $\varphi(1) - \varphi(0)$ with change in Reynolds number, R_{θ} .

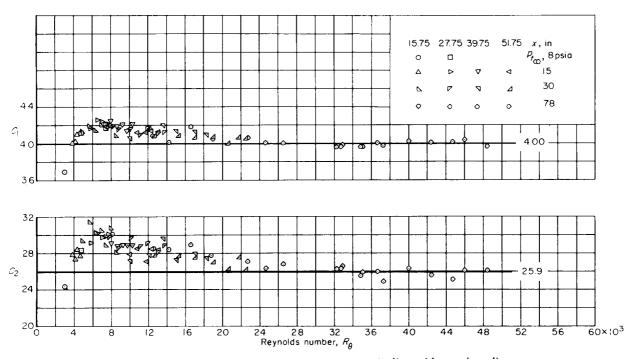
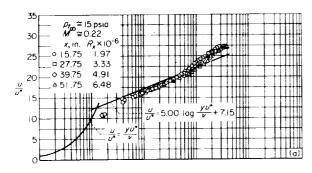


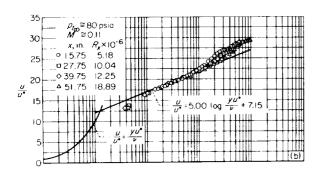
Figure 17. -Variation of C_1 and C_2 with change in Reynolds number, R_{θ} .

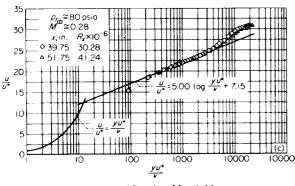
 $509374-59 \ \cdot \ 3$

is dependent principally on C_1 and C_2 , both of which change markedly below Reynolds numbers of about 21 million from the asymptotic value used in the computation (see fig. 17).

Average and local skin-friction coefficient.—It appears in figures 12 and 13 that the computed values of both the average skin-friction coefficient and the local skin-friction coefficient represent the measured values quite well for Reynolds numbers as low as about five or six million.—The scatter in the data presented for both the average and local



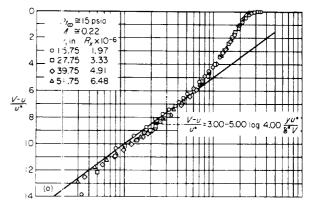


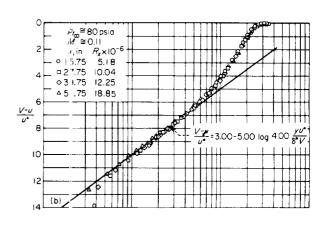


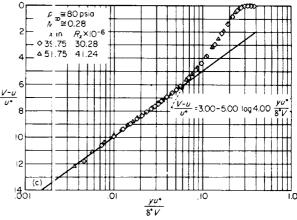
- (a) $p_{I\omega}=15$ psia, $M\approx 0.22$
- (b) $p_{t\infty} = 80 \text{ psia}, M = 0.11$
- (c) $p_{t\infty}$ =80 psia, $M\!=\!0.28$

Figure 18. Boundary-layer velocity profiles in terms of the "wall law."

skin-friction coefficients is represented generally by a change of skin friction of about ± 1 percent, and this s also about the variation of the measured







- (a) $p_{i\omega} = 15$ psia, M = 0.22
- (b) $p_{t\infty} = 80 \text{ psia}, M = 0.11$
- (e) $p_{t_{\infty}} = 80$ psia, M = 0.28

Figure 19.—Boundary-layer velocity profiles in terms of the "velocity-defect law."

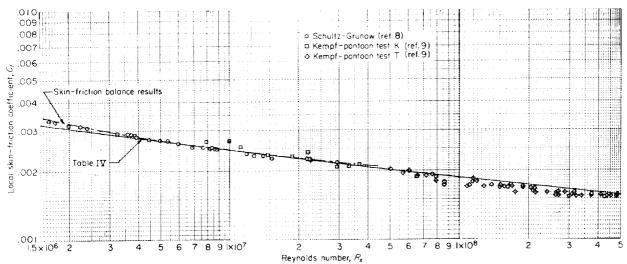


FIGURE 20.—Comparison of previous existing data with the measured data and with the curve computed using the constants derived from the present data.

from the computed skin-friction coefficient at higher Reynolds numbers.

Law of the wall and velocity-defect law. In figures 18 and 19 it is shown that the wall law and velocity-defect law derived using the value of the constants in the range of Reynolds number independence do not represent the measured data except at the higher values of Reynolds number. Here again this is easily understood after inspection of the variation with Reynolds number, shown in figure 16, of the parameters used in both laws.

COMPARISON OF MEASURED DATA AND COMPUTED FRICTION LAW WITH MEASURED DATA OF OTHER INVESTIGATIONS

Local skin-friction coefficients.-There are presented in figure 20 the local skin-friction coefficients measured by Schultz-Grunow (ref. 8) in an air channel and by Kempf (ref. 9) on a pontoon in water. In the region of Reynolds number where the two sets of data overlap, Kempf's data appear to be somewhat higher than those of Schultz-Grunow. In this region of overlap, Schultz-Grunow's data agree quite well with the skinfriction balance results. Comparison of the measured local skin-friction-coefficient data of the present investigation with those of both Schultz-Grunow and Kempf indicates remarkable agreement in the Reynolds number range of the investigation when it is considered that the data came from three grossly different pieces of equipment. Here, as in figure 13, there is a tendency for the measured data to be higher than the computed friction law (table IV) for Reynolds numbers smaller than about 4 or 5 million. However, the computed friction law does match the measured data quite well for a range of Reynolds numbers from 4 or 5 million to about 60 million. For Reynolds numbers above 60 million Kempf's data appear to fall below the line representing the computed law.

Average skin-friction coefficient. —A comparison of the measured average skin-friction coefficients of the present investigation with the Schoenherr line (ref. 11) is presented in figure 12. The Schoenherr line gives larger values of skin friction than were measured in the present investigation for Reynolds numbers from 3 to 30 million, but became equal to the measured values at Reynolds number from 1 to 3 million and from 30 to 45 million. The measured data are best represented by the Schoenherr line in the range of Reynolds numbers from 1 to 3 million and by the computed law (table IV) in the range of Reynolds numbers from 5 to 45 million.

Figure 21 is a reproduction of a figure presented in reference 11 with the exception that the computed friction law of the present investigation is also presented for comparison. The friction law as computed from the data of the present investigation gives values of skin-friction coefficient as much as 8 percent lower than the Schoenherr line at a Reynolds number of 1 million and as much as 6 percent higher at a Reynolds number of 1×10^9 . Similar to previous comparisons be-

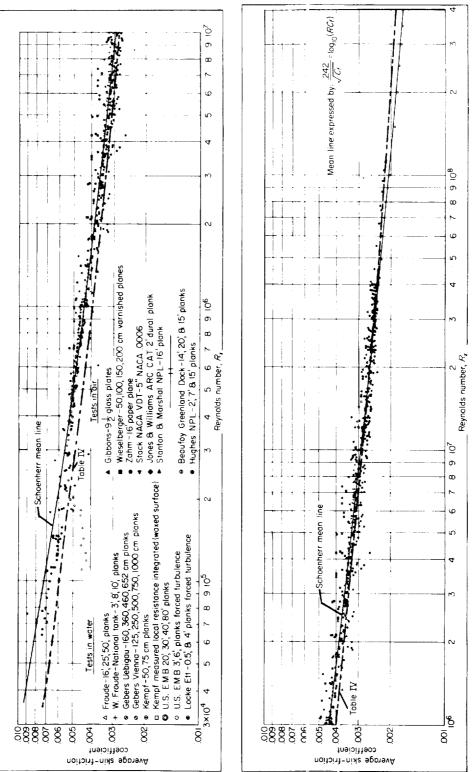


FIGURE 21.—Comparison of previous existing data with computed friction law (table IV).

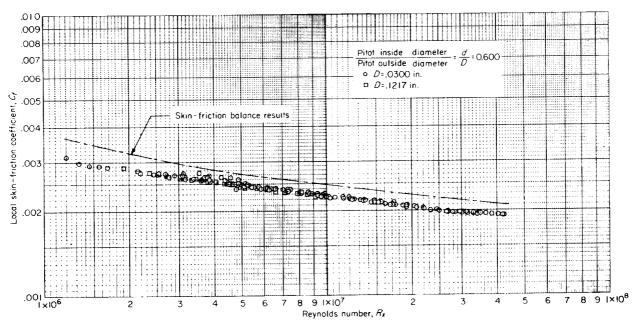


Figure 22. Variation of local skin-friction coefficient with change in Reynolds number, R_x ; Preston tube technique.

tween the computed curve and the measured data, the measured data are somewhat higher in the low Reynolds number range. In the range of Reynolds number from 5 million to 100 million the computed law seems to represent the data quite well. Beyond a Reynolds number of 100 million there is only one set of data available to compare with the computed values and they lie below the computed line for all higher values of Reynolds number.

A SIMPLE METHOD FOR DETERMINING LOCAL SURFACE-SHEAR STRESS IN A TURBULENT BOUNDARY LAYER

There are presented in figure 22 the results of measurements of local skin-friction coefficient using a calibrated total-head tube as proposed by J. H. Preston in 1953 and previously described in the section on experimental methods. On the same figure is presented a line representing the faired value of the data measured with the floating-element device as presented previously in figure 13. In general, the Preston tube device indicates a smaller skin friction than the floating-element device. However, the results of both methods can be made to agree quite well if the calibration presented by Preston in reference 5 is modified slightly.

From the work of Preston it has been shown that the calibration of the tubes is valid only if the value of the expression $\log_{10} \frac{(p_t - p)d^2}{4 - r^2}$ is

greater than about 5.0 but less than about 7.5. These limiting values also seem to be the limiting values obtained in the present investigation. When the value of the expression $\log_{10} \frac{(p_t - p)d^2}{4\rho \nu^2}$ falls outside of these limits the measured skin friction immediately varies away from the general trend of similar data measured at the same Reynolds number when the value of the logarithmic expression falls within the prescribed values.

It appears that the Preston tube device can be quite useful in measuring the local surface-shear stress in a turbulent boundary layer where the longitudinal static-pressure gradient is zero. Not only does it appear to be accurate but it is extremely simple and inexpensive to construct. Also, the indicating equipment is simple and readily available to most investigators.

For Reynolds numbers greater than 2.5 million the revised calibration suggested by the measured surface-shear stress data obtained on the floatingelement device is

$$\log_{10} \frac{\tau_u d^2}{4 \rho v^2} = -1.366 + 0.877 \log_{10} \frac{(p_t - p) d^2}{4 \rho v^2}$$

as compared with Preston's calibration of reference 5,

$$\log_{10} \frac{\tau_u d^2}{4\rho \nu^2} = -1.396 + 0.875 \log_{10} \frac{(p_t - p)d^2}{4\rho \nu^2}$$

For Reynolds numbers lower than 2.5 million use

of the revised calibration results in values of surface shear which are lower than the measured data.

CONCLUDING REMARKS

The measured local skin-friction coefficients obtained from the floating-element skin-friction balance agree well with the long accepted experimental data of Schultz-Grunow and Kempf in the range of Reynolds numbers from about 1 million to about 45 million.

The average skin-friction coefficients deduced from the measured velocity profiles are generally below the Schoenherr line except at the lowest values of Reynolds number. As the Reynolds number approached 45 million, the highest value attained in the present investigation, the measured average skin friction became equal to the value predicted by Schoenherr. However, the rate of change of the measured average skin-friction coefficient with increasing Reynolds number is smaller than that predicted by Schoenherr.

The frictional drag experienced by a flat-plate surface has been computed by both the momentum-defect method and the integration of the local surface shear. At values of Reynolds number from 14 million to 45 million the results of both methods are in good agreement but show a discrepancy of as much as 4 percent in the range of Reynolds numbers from 2 to 6 million.

In the light of the data of the present investigation a new frictional resistance law for a smooth plate having zero pressure gradient may be written. However, there is some doubt as to the absolute values of the experimentally determined parameters which must be used in conjunction with the skin-friction equations to write a law. These parameters appear to approach asymptotically a constant value, as was anticipated. As a result of the scatter in the values of the parameters obtained from the experimental data and the limited Reynolds number range attained in the investigation, there seems to be some doubt as to the validity of a law written on the basis of these parameters.

The local skin friction determined from measurements utilizing a calibrated pitot tube mounted on the surface as proposed by J. H. Preston had a lower value than that measured by the floating-element skin-friction balance. However, a small adjustment of Preston's calibration of the pitot tube brought the two results into good agreement.

The Preston pitot tube appears to be an inexpensive and accurate device for making local surface-shear-stress measurements.

AMES RESEARCH CENTER

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION MOFZETT FIELD, CALIF., Dec. 9, 1957

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TABLE L-MEASURED BOUNDARY-LAYER VELOCITY PROFILES

(a) x = 15.75 inches

£5883 28.83 28.65 23 0.187 y, inch 0283 0.0335 0.0335 0.0335 0.0335 0.0335 0.0335 1935 1435 1635 1835 医胃炎医胃 生物医疗 医肾髓管 医胃炎 2,65 28,41 69 0,152 0535 0735 0735 0935 おき置置者 25.26 25.26 2.367 2.367 y, inch 0535 0535 0735 0785 0835 0285 0385 0485 0485 0485 1135 1135 1135 1135 1635 12.21 14.84 0.20 84.20 0285 0335 0435 0435 0435 0435 0635 0735 0735 0835 0835 충돌물림 \$5.88 8.88 8.88 14.83 15.83 0.24 0.0035 0085 0135 0135 0135 0235 0335 0435 0485 0535 0735 0735 0835 0835 1035 1235 1435 1835 1835 0 41 683 683 671 671 775 775 786 786 787 787 787 787 787 787 86888 25.25 27.81 0.21 0.0035 .0085 .0135 .0185 .0235 0285 0385 0485 0485 .0535 .0635 .0735 .0935 1035 1135 1635 1835 8221224 | 8824 | 8821228 8221224 | 8824 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 8825 | 88 28232 25 26232 25 1.97 15.17 71 0.215 0.0035 .0085 .0135 .0185 2035 2435 2435 2635 2835 3035 0535 0635 0735 0835 0935 1035 1235 1435 1635 1835 76 0.308 \$21288 **\$3228** \$25.58 15.58 936 952 979 979 987 8 8 8 8 8 $\begin{array}{c} - \\ R_L \times 10^{-6} \\ P_{CR} \text{(psia)} \\ T \text{ (°F)} \\ M \end{array}$ $\begin{array}{c} 0.0035 \\ .0085 \\ .0135 \\ .0175 \\ .0235 \end{array}$ 0535 0635 0735 0835 0935

TABLE L.—MEASURED BOUNDARY-LAYER VELOCITY PROFILES—Continued

TABLE I.-MEASURED BOUNDARY-LAYER VELOCITY PROFILES-Continued

(b) r = 27.75 inches

ALC &	1/1	0.461 .571 .611 .635	668 693 714	843 1717 1761 1792 1778 1778 1778 1778 1778 1778 1778 177	889 889 909 927 843	980 983 991 385	866 666 666	
5,82 29,45 76 0,186	y, inch	0.0035 .0085 .0135 .0135	. 0285 . 0335 . 0435 . 0485	. 0535 . 0785 . 1035 . 1285 . 1535	. 1785 . 2035 . 2285 . 2536 . 2785	. 3035 . 3285 . 3535 . 3785 . 4035	. 4285 . 4535 . 4785	
-1 25	n/1'	0, 451 564 601 629 647	690 690 70 709 709	7.7. 7.7. 7.87. 839	88.4 905 923 940 940	956 988 988 989 989	706 388 886 866	
4, 90 29, 51 65 0, 152	y, inch	0.0035 ,0085 ,0135 ,0135	0285 0335 0435 0435		2285 2285 2585 2785 2785	3035 3285 3535 3785 4035	. 4285 . 4485 . 4785 . 5035	
6 4 00	u/1'	0. 463 . 555 . 602 . 627 . 643	26588 26788 2678 2678 2678 2678 2678 2678 2	28. 28. 28. 28. 28. 28. 28. 28. 28. 28.	859 920 938 938	954 967 979 987 992	995 998 998 998	866 866
83 94 0.310	y, inch	0. 0035 . 0085 . 0135 . 0185	0285 0335 0435 0435 0485	. 0535 . 0785 . 1035 . 1285 . 1535	1785 2035 2285 2535 2785	3035 3285 3535 3785 4035	. 4285 . 4535 . 5035 . 5285	. 5535
କଳ ୬	1/1/1	0. +46 560 601 624 645	658 672 693 703	E1-15-18-28-28-28-28-28-28-28-28-28-28-28-28-28	880 880 900 918 935	952 967 977 988 988	668 668 668 668 668 668 668 668 668 668	866
2.50 15.01 0.288	y, inch	0.0035 .0085 .0135 .0185	0285 0385 0438 0438		1785 2035 2285 2585 2585 2785	3035 3285 3535 3785 4035	4285 4535 4785 5035 5285	. 5535
<u>.</u>	n/1.	0, 438 558 599 617 637	653 657 691 699	2005 805 805 805 805 805 805 805 805 805	854 875 895 915 934	949 974 974 984 991	999 2008 8008 8008	
4, 17 15, 07 78 0, 266	y, inch	0.0035 .0095 .0135 .0185 .0235	.0285 .0335 .0385 .0435	0635 0785 1035 1285 1535	2035 2283 2585 2585 2785	3035 3285 3535 3785 4035	4285 4535 4785 5035	
	n/1.	0. 433 .554 .588 .615 .638	655 655 655 655 655 656 656	25.E. 25.25.25.25.25.25.25.25.25.25.25.25.25.2	850 871 1929 1929	946 946 971 982 990	266 666 666 666	
3.83 15.01 74 0.242	y, inch	0. 0035 . 0085 . 0135 . 0185	0285 0335 0435 0485	0535 0785 1035 1285 1535	2285 2585 2585 2585 2785	3025 3285 3585 3786 4035	4285 4535 4785 5035	
38 24 24	11/11	0. 428 . 537 . 584 . 613	659 672 685 694	702 771 771 798 826	849 871 890 909 927	986 980 980 980	998 998 998 998	666
3, 33 14, 75 76 0, 214	y, inch	0, 0035 , 0085 , 0135 , 0185	0.02% 0.03% 0.03% 0.04%	0535 0785 11035 1285 1535	1785 2035 2285 2535 2785	3035 3285 3535 3785 4035	4285 4535 4785 5035 5285	. 5535 . 5785
68-12 0.309	11/11	0, 391 533 580 599 617	. 15.69 1.69 1.69 1.69 1.69 1.69 1.69 1.69 1	687 723 781 805 805	828 851 894 894	926 942 956 968 977	986 991 994 996 997	998 998 888 888
$P_{t,\infty}(\text{psia})$ $T_{t,\infty}(\text{PSi})$ $T_{t,\infty}(\text{PS})$	y, inch	0.0035 0.0085 0135 0185 0235	.0285 .0335 .0385 .0435	.0535 .0785 .1035 .1285 .1535	1785 2035 2285 2535 2785	3035 3285 3535 3785 4035	4285 4535 4785 5035 5285	. 5535 . 5785 . 6035 . 6285

TABLE L-MEASURED BOUNDARY-LAYER VELOCITY PROFILES -Continued

		1//n	0. 492 . 603 . 649 . 673 . 683	25.52.52 25.52.52 25.52.52 25.52.52 25.52.52 25.52.52 25.52.52 25.52.52 25.52	25.59 26.20 25.00	891 947 963 963	986 986 990 998	666 - 666 -	
	14. 08 84. 93 67 0. 152	y, inch	0, 0035 , 0085 , 0135 , 0185 , 0235	0285 0335 0335 0435 0435	0535 0785 1035 1285 1535	1785 1785 2882 2885 2885 2885 2785	3035 3285 3535 3785 4035	42N5 4535 4785	
.	#8 E	11/11	0. 482 . 591 . 636 . 659 . 678	985 7119 727 727 727 727	75.55 55.55	880 921 936 955	986 888 888 888 888 888	866 866	
	10.04 85.28 66 0.108	y, inch	0.0035 .0085 .0135 .0135 .0185	0285 0335 0385 0435 0485	0535 0785 1035 1285 1585	2285 2285 2285 2585 2585 2785	3025 3285 3535 3785 4035	4285 4535 4785	 ! ! !
	33.0 310	и/11	0.490 .596 .635 .661	696 708 719 729 738	\$2.2 \$2.0 \$3.0 \$3.0 \$3.0 \$3.0 \$3.0 \$3.0 \$3.0 \$3	988 921 938 938 938	970 988 988 499 799	866 666 666 666	666
lded	9.83 30.83 65.63 01.810	y, inch	0.0035 .0085 .0135 .0185	0285 0335 0435 0435		1785 2285 2285 2585 2785 2785	3035 3285 3535 3785 4035	4285 4535 4785 5035 5285	. 5535
(b) x=27.75 inches—Conciuded	25. 32.	n/1.	0, 479 595 629 655 676	25 H 75 8	######################################	879 921 938 938	969 979 988 998 998	888 886 886 886 886 886 886 886 886	1
(b) $x = 27.75$	9, 00 30, 32 78 0, 28 0, 28	y, inch	0.0035 .0085 .0135 .0185	0285 0335 0435 0435	0535 0785 11035 1285 1535	1785 2285 2585 2785 2785	3035 3285 3535 3785 4035	4285 4535 4785 5035	
	25.27 13.62 1.265 1.265	11/11	0. 483 . 587 . 625 . 653	788 227 228 228 228 228 228 228 228 228	£15,555	878 900 917 936 953	9.95 87.89 9.99 9.99 9.99	886 866 866	
	∞ इत्तंह⊂	y, inch	0.0035 .0085 .0135 .0185	0285 0385 0485 0485		1785 2285 2535 2785	3035 3285 3535 3785 4035	. 4285 . 4535 . 4785 . 5085	
	7. 58 29.10 68 0. 243	11/11	0. 473 . 579 . 619 . 646	388 115 115 115 115 115 115 115 115 115 1	#E##	85 9 9 7 7 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	9.83 9.83 9.83 9.83 9.83 9.83 9.83	6886 6896 6800	
	- SI & T	y, inch	0,0035 .0085 .0135 .0135	.0285 .0335 .0385 .0435	.0535 .0785 .1035 .1285 .1535	1785 2285 2585 2535 2785	3035 3285 3535 3785 4035	22. 25. 25. 25. 25. 25. 25. 25. 25. 25.	
	29, 46 29, 46 74 0. 217	1/11	1. 469 1. 578 1. 644 1. 644	85 58 57 58 58 57 57 58 58 58 58 58 58 58 58 58 58 58 58 58 5	£8333	913 932 948	965 975 988 988 988	866	
	$\begin{array}{c} R_x \times 10^{-6} \\ p_{t,\infty} \text{ (bsia)} \\ T({}^{\circ}\text{F}) \\ M \end{array}$	y, inch	0.0035 .0085 .0135 .0185	.0285 .0335 .0385 .0435	. 0535 . 0785 . 1035 . 1285 . 1535	. 1785 . 2035 . 2285 . 2535 . 2785	3035 3285 3535 3785 4035	288 288 388 388	

TABLE L-MEASURED BOUNDARY-LAYER VELOCITY PROFILES—Continued

(c) x = 39.75 inches

88 H	,1/n	0. 445 . 558 . 599 . 626	. 661 674 685 . 703	712 744 768 791 810	828 846 879 836 936	957 976 989 995	666			
9.56 29.26 77 0.217	y, inch	0.0035 .0085 .0135 .0185 .0235	.0285 .0335 .0385 .0435	. 0535 . 0785 . 1035 . 1285 . 1535	. 1785 . 2035 . 2535 3035 . 3535	. 4035 . 4535 . 5035 . 5535 . 6035	. 6535	!	:	
31 33 187	n/1.	0. 439 . 564 . 620 . 623 . 643	. 657 . 689 . 689 . 689	708 785 785 785 708	825 876 905 831	9999999				
8.31 29.03 73 0.187	y, inch	0.0085 .0085 .0135 .0185	.0285 .0335 .0385 .0435 .0485	0535 0785 1035 1285 1535	1785 2035 2535 3035 3535	4035 4535 5035 5535 6035				
2 75 E	n/1	0. 431 . 555 . 590 . 613	645 658 672 682 690	697 730 782 782 801	836 836 871 899 899	. 952 971 985 994 998				
6.74 28.72 69 0.151	y, inch	0.0035 0035 0135 0185	0285 0335 0385 0435 0485	0535 0785 1035 1285 1535	1785 2035 2535 3035 3535	4035 4535 5035 5535 6035				
85 90 310	n/11	0. 434 560 593 618 635	# 6 6 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	1325 1325 148 148	7.77.7. 7.77.7. 7.02. 7.03. 7.03. 828.	8.836 8.836 8.77 8.836 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	922 933 942 942	. 951 . 960 . 968 . 975 . 980	985 993 995 995	866 866 866
6,78 14.90 79 0.310	y, inch	0.0035 .0085 .0135 .0185	0285 0335 0435 0435	0535 0735 0735 0835 0935	1035 1235 1435 1635 1835	2035 2235 2435 2435 2635 2835	. 3035 . 3235 . 3435 . 3635 . 3835	4035 4235 4435 4635 4835	5035 5235 5435 5635 5635	6035 6235 6435 6635
36 89 288	n/11	0. 426 554 591 613 631	749 669 670 670 888		818 834 898 898 898 927	950 971 985 994 998	. 999 1. 000 1. 000			!
6.36 14.89 77 0.288	y, inch	0.0035 .0085 .0135 .0185	. 0285 . 0335 . 0435 . 0485	. 0535 . 0785 . 1035 . 1285	. 1785 . 2035 . 2535 . 3035 . 3535	. 4035 . 4535 . 5035 . 5535 . 6035	. 6535 . 7035 . 7535	: !	! !	
5.96 14.87 73 0.267	1/1/	0, 420 . 552 . 587 . 612 . 629	645 655 675 675	692 726 750 773 794	812 831 865 896 923	947 983 998 998	666 686		I	
44 44 EC 9.	y, inch	0.0035 .0085 .0135 .0185	. 0285 . 0335 . 0435 . 0435	. 0535 . 0785 . 1035 . 1285 . 1535		. 4035 . 4535 . 5035 . 5535	. 6835 . 7035	i :	· 	: : : !
5.55 14.79 64 0.243	n/17	0. 414 . 547 . 583 . 606 . 606	7.83 7.83 6.50 1.79 1.79	687 747 769 790	808 828 828 861 891 919	24.88.88.88.88.88.88.88.88.88.88.88.88.88	33333			
4 H 20	y, inch	0.0035 .0085 .0135 .0185 .0235	. 0285 . 0335 . 0335 . 0435 . 0485	. 0535 . 0785 . 1035 . 1285	. 1785 . 2035 . 2535 . 3035	. 4035 . 4535 . 5035 . 5535	6535 7035 7535 8035 8535			!
4.91 14.79 71 0.217	11/11	0.401 .542 .578 .599	646 655 672	681 715 743 784	803 822 857 887 915	941 960 990 990	866 666 666			
$R_x \times 10^{-6}$ $p_{t_\infty}(\text{psla})$ $T(^{\circ}\text{F})$ M	y, inch	0.0035 .0085 .0135 .0185	0285 0335 0435 0435	.0535 .0785 .1035 .1285	. 1785 . 2035 . 2536 . 3035	. 4035 . 4535 . 5035 . 5535	. 6535 . 7035 . 7535			

TABLE I.—MEASURED BOUNDARY-LAYER VELOCITY PROFILES.--Continued

(c) x = 39.75 inches -Continued

1 1 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	1//n	0. 493 . 603 . 645 . 669	 15.57.68 15.57.68 15.57.68	85 G 85	836 836 836 854 854	900 928 951 970 985	99. 99. 99. 998. 998.
26.21 78.34 100 0.235	y, inch	0.0035 .0085 .0135 .0185	. 0285 . 0335 . 0435 . 0435	.0535 .0535 .0735 .0835	1035 1285 1535 1785 2035	2535 3035 3535 4035	. 5035 . 5535 . 6035 . 6535
	1/n	0. 488 603 645 668 686	85 85 85 85 85 85 85	966668 966668	797 783 835 853 870	900 927 952 970 985	988 888 987 758 888 888
25.82 78.80 0.231	y, inch	0.0035 .0085 .0135 .0185	0285 0335 0435 0485	0525 0635 0735 0835 0935	1035 1286 1535 1785 2035	. 2535 . 3035 . 4035 . 4535	. 5035 . 5535 . 6035 . 6535 . 7035 . 7335
30 X	1/1/n	0.520 . 610 . 645 . 668	88. 95. 55. 80. 52. 55. 82. 55.	25.25.25.25.25.25.25.25.25.25.25.25.25.2	. 796 . 815 . 834 . 852 . 852	886 926 949 971 986	. 995 . 999 1. 000 1. 000
21. 29 75.87 81 0. 187	y, inch	0.0035 .0085 .0135 .0185	. 0285 . 0335 . 0385 . 0435 . 0485	0535 0635 0735 0835 0935	1035 1235 1735 2035	2535 3035 3535 4035 4535	. 6535 . 6535 . 6535
34 57 152	.1/n		51.52.52	766 789 829 829 847	8.8.2.8.8.8. 8.8.2.8.8.8.	886 886 886 886 886 886	
17.34 74.57 75.00 0.152	y, inch	0.0035 .0085 .0135 .0185	.0335 .0385 .0435 .0485 .0585	0785 1035 1285 1535 1785	. 2035 . 2535 . 3035 . 3535 . 4035	. 4535 . 5535 . 5535 . 6035	
	u/1.	0. 455 568 612 638 658	25. 26. 26. 27. 26. 27. 26. 27. 27. 27. 27. 27. 27. 27. 27. 27. 27	722 7754 778 817	25.55.25.25.25.25.25.25.25.25.25.25.25.2	941 980 980 998 998	1.000
12, 25 73, 39 69 69 0, 108	y, inch	0.0035 .0085 .0135 .0185	.0335 0335 0385 0485 0485	.0535 .0785 .1035 .1285	1785 2285 2585 3035 3035	. 3535 . 4035 . 4535 . 5035	. 6535
55 311	11/11	0. 460 . 575 . 614 . 642	858 805 805 717	55.55 55.55	85.5 87.8 98.9 91.6 91.6	941 986 997 768	1.000
13.45 30.55 94 0.311	y, inch	0.60035 .00085 .0135 .0185	. 02885 . 03885 . 04885 . 04885	.0535 .0785 .1035 .1285 .1585	.1785 2286 2535 3035	. 3535 . 4035 . 4535 . 6035	. 6535
88 88 80 80 80 80 80 80 80 80 80 80 80 80 8	11/11	0. 455 . 570 . 613 . 639	. 654 . 686 . 686 . 707 . 715	227. 887. 818.	886 914 940 940	9.99.97.99.99.99.99.99.99.99.99.99.99.99	986
12, 53 30, 35 93 93 0, 290	y, inch	0. 0035 . 0085 . 0135 . 0185 . 0235	.0285 .0335 .0385 .0485	0535 0785 1035 1285 1535	1785 2535 3635 3635	. 4035 . 4535 . 5035 . 6035 . 6035	. 6535
10.68 29,75 85 0.244	1/n	0. 450 562 632 . 652	. 667 . 679 . 691 . 709	747 747 795 814 814	288 288 198 1988	9.50 8.00 8.00 8.00 8.00 8.00 8.00 8.00 8	686
$P_{t_{\infty}}(psia)$ $T_{(^{\circ}F)}$ M	y, inch	0.0035 .0085 .0135 .0185 .0235	. 0235 . 0335 . 0435 . 0485	. 0535 . 0785 . 1035 . 1285 . 1535	1786 2503 2533 3038 3538 3538	. 4035 . 4535 . 5035 . 5535 . 6035	. 6585

TABLE I. MEASURED BOUNDARY-LAYER VELOCITY PROFILES -Continued

(e) x=39.75 inches Concluded

$P_x \times 10^{-6}$ $p_{t_{\infty}}$ (psia) T (°F) M	30.29 80.51 107 0.272	28.1 79.1 108		28.11 78.22 90 0.247			
M	0.272	0.	200	0.			
y, inch	u/V	y, inch	u/V	y, inch	u/V		
0, 0035	0. 509	0, 0035	0.511	0.0035	0, 488		
. 0085	. 614		. 614	. 0085	, 609		
. 0135	653	, 0135	. 652	. 0135	648		
, 0185	676	, 0185	. 675	. 0185	. 671		
0235	. 694	. 0235	.692	. 0235	. 690		
, 0285	. 708	0285	. 705	. 0285	. 704		
. 0335	. 720	. 0335	. 718	. 0335	. 716		
. 0385	. 728	. 0385	. 727	. 0385	. 725		
. 0435	. 738	. 0435	. 736	. 0435	734		
. 0485	. 744	. 0485	744	. 0485	. 741		
, 0535	. 751	. 0535	. 751	. 0535	. 749		
. 0585	. 756	. 0635	. 763	. 0635	. 760		
. 0635	. 763	. 0735	. 773	. 0735	. 771		
. 0735	. 774	0835	. 783	. 0835	. 781		
. 0835	. 784	. 0935	. 792	0935	. 791		
. 0935	. 792	. 1035	. 801	. 1035	. 799		
. 1035	. 801	. 1285	. 821	. 1285	. 820		
1285	. 821	. 1535	840	, 1535	839		
. 1535	. 841	1785	. 858	. 1785	. 855		
. 1785	858	. 2035	. 873	2035	872		
. 2035	. 874	. 2535	904	. 2535	902 . 929		
. 2535	. 903	3035	. 930	. 3035	. 929		
. 3035	. 931	3535	. 956	. 3535	. 972		
. 3535	. 955	. 4035	. 974	. 4535	. 986		
4035	. 973	4535	. 987	4000			
. 4535	. 987	5035	. 996	. 5035	. 993		
. 5035	. 995	. 5535	. 999	. 5535	, 997		
. 5535	. 998	6035	1,000	. 6035	. 998		
. 6035	, 999	. 6535	1,000	. 6535	. 998		
. 6535	. 999	. 7035	1,000	7035	. 998		
. 7035	. 999	7535	1,000	. 7535	. 998		
. 7535	, 999	. 8035	1,000	1			
, 8035	. 999			İ			
. 8535	, 999						
. 9035	. 999						

TABLE L.-MEASURED BOUNDARY-LAYER VELOCITY PROFILES—Continued

	1/n	888=8	\$288c\$	25 248 248 288 288 288	26.25.4		: : ! ආජුවුවල	 :	
10.95 29.81 0.190	/n	0.430 .553 .589 .611 .630	645 658 677 686 686	**************************************	*2020	18.88.88.88.88.88.88.88.88.88.88.88.88.8	33333		
= 88.50	y, inch	0.0035 .0085 .0135 .0185 .0235	0285 0335 0385 0435	0535 0785 1035 1285 1535	1785 2035 2535 3035 3535	4035 4535 5035 5535 6035	6535 7635 8035 8535 8535		
8, 96 29, 66 83 0, 155	1/11	0, 428 , 538 , 604 , 604	\$4868	387.78	797 8 8 8 8 12 1 1 1 1 1 1 1 1 1 1 1 1 1 1	25 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	. 986 . 985 1. 000	:	
x ည်ညွှင် 	y, inch	0.6035 .0085 .0185 .0185	0285 0335 0435 0435	.0535 .0785 .1035 .1285 .1535	1785 2035 2535 3035 3535	4035 4537 5035 5535 6035	. 6535 . 7035 . 7535		
05 20 317	1/1/	0. 409 543 579 602 620	635 647 658 669 678	25 24 25 28 28 28	25. 25. 25. 25. 25. 25. 25. 25. 25. 25.	888 888 888 888 888 888 910	920 931 941 950	966 974 983 988	666 666 666 666 666 666 666 666 666 66
9.05 15.20 85 0.317	y, inch	0.0085 .0085 .0135 .0185	0285 0335 0385 0435 0485	0535 0785 1035 1285 1535	1785 2035 2285 2585 2585 2785	3035 23035 3535 3785 4035	4285 4785 5035 5285	5535 5785 6435 6285 6535	525 525 525 525 525 525 525 525 525 525
52 13 295	n/Γ	0. 405 540 576 600 6115	644 657 656 673 673	184 184 184 184 184 184 184 184 184 184	200 200 200 200 200 200 200 200 200 200	909 920 939 949 958	965 972 977 984 988	9922 994 998 998 999	
8. 52 15. 13 81 0. 295	y, inch	0, 0035 (0085 (0185 (0185	0335 0335 0435 0485	0535 0785 1035 1285 1535	1785 2035 2535 3035 3535	4035 4525 4785 5035 5285	5535 5785 6735 6235 6535	6785 7085 7385 7385 7385 7385	
92 90 90 974	1/11	0.398 535 572 598 616	643 643 662 663 663 672	680 711 734 756 775 775	25. 25. 25. 25. 25. 25. 25. 26. 26. 26. 26. 26. 26. 26. 26. 26. 26	918 956 976 982	991 998 999 999		i
7. 92 14. 90 76 0. 274	y, inch	0. 0035 . 0085 . 0135 . 0185	0285 0335 0385 0485 0485	0535 0785 1035 1285 1585	2285 2785 3285 3285 3785	4285 4785 5285 5785 6285	6785 7785 8285 8785 8785		
55 SE	1/1/	0.386 .527 .568 .591 .609	624 645 645 687 688	674 707 732 751 751 768	800 832 885 885 905	925 045 941 973 986	. 991 . 997 . 999 . 999 . 1. 000		
. 1. 15 . 2. 248 . 2. 248	y, inch	0.0035 0085 0135 0185	0285 0335 0385 0435 0485	. 0535 . 0785 . 1035 . 1285 . 1535	2035 2535 3035 3535 4035	. 4535 5045 . 5535 . 6035	.7085 .7585 .8035 .8535 .9035		
**************************************	1/11	0. 390 524 506 586 604	619 633 644 653 653 653	845 845 845 845 845 845 845 845 845 845	853 877 878 878 878	. 942 958 958 976	983 998 998 1,000		-
6. 44. 44. 44. 44. 44. 44. 44. 44. 44. 4	y, inch	0.0035 .0085 .0135 .0185	0335 0335 0335 0435 0435	0535 0785 1035 1285 1535	. 1785 . 2035 . 2535 . 3035 . 3535	. 4035 4735 5035 . 5535 . 6035	. 6385 . 7085 . 7585 . 8685 . 8585	 	
4. 78. 7. 97. 81. 0. 317	.1/n	0.336 . 497 . 542 . 569 . 589	623 633 640	648 680 705 725 745	761 883 883 883 883	885 927 945 945	9.56 9.69 9.69 9.69 9.69 9.69	998 998 998 1.000	
$P_{r_{\infty}} (\text{psia})$ $T^{(\circ)}(F)$	y, inch	0. 0035 . 0085 . 0135 . 0185	.0285 .0335 .0335 .0435 .0435	0535 0785 1035 1285 1535	1785 2035 2535 3035 3535	4035 4 535 5035 5535 6035	6535 7035 7535 8035 8535	. 9035 9535 1. 0035 1. 0535	

TABLE I. -MEASURED BOUNDARY-LAYER VELOUITY PROFILES-Continued

(d) x = 51.75 inches Continued

30. 55 80. 49 77 0. 193	1/n	0,457 589 . 631 . 656	.688 .709 .709 .718 .724	731 780 781 800 815	830 844 871 885 816	983	. 996 . 999 1. 000 1. 000	
8815	y, inch	0.0035 .0085 .0135 .0185	.0285 .0335 .0385 .0435 .0485	0535 0785 1035 1285 1535	2035 2035 2535 2535 3035 3535	. 4035 . 4535 . 5035 . 5535 . 6035	. 6535 7035 7735 8035 8035 . 8535	
15 to	11/11	0. 458 . 582 . 624 . 652	683 695 705 720	121 121 128 138 138 138	28.88.88.89.89.89.813.89.89.89.89.89.89.89.89.89.89.89.89.89.	980 980 989	. 996 . 998 . 999 1. 000 1. 000	1.000
26, 75 85, 15 71 0, 157	y, inch	0, 0035 . 0085 . 0135 . 0185	0285 0335 0435 0435	. 0535 . 0785 . 1035 . 1285	2035 2035 3035 3535 3535	. 4035 . 4535 . 5035 . 6035	. 6535 . 7035 . 7535 . 8035 . 8535	. 9035 . 9535
=	11/11	632 635 655 655	689 170 108 108 108 108 108	27.7.7.2.8. 20.2.7.7.7.8.8.8.8.8.8.8.8.8.8.8.8.8.8.8.8	857 857 882 903	. 924 . 941 . 961 . 955	206 308 308 308 308	
18.89 84.01 67 0.111	y, inch	0.0035 .0085 .0135 .0185	. 0285 . 0335 . 0435 . 0435		2035 2035 2535 3035 3535	4035 4535 5035 5535 6035	. 6535 . 7035 . 7535 . 8035	
# L1	11/11	0. 429 . 466 . 575 . 612	25.00 25.00	705 783 161 197 797	8.8212 8.8212 8.7249 8.02	923 941 943 943	. 992 . 999 . 999 . 999	1.000
2.82.82 2.82.82 2.82 2.82 2.82 3.82 3.82	y, inch	0,0035 .0085 .0135 .0185	0285 0335 0385 0435	0635 0785 1035 1285 1635	2035 2035 2535 3035 3535	4035 4535 5035 6035	.6535 .7635 .7535 .8035 .8535	. 9035 . 9535 1. 0035
# 3 Q:	11/11	0, 458 573 609 634 632	665 678 788 705 705	217.7 7.7.7 7.82 808 809	828 828 882 882 882 883	25.25.25.25.25.25.25.25.25.25.25.25.25.2	993 998 998 998	
16.61 22.76 38. 36. 0.295	y, inch	0.0035 .0085 .0135 .0185	028 0335 0435 0435 0485	0535 0785 1035 1285 1535	2035 2035 2535 3035 3535	4635 5635 5635 6035	5535 7035 7535 8035 8535	: :
6 73	n/17	0, 450 569 607 631 648	88.67.88.49.55 10.00	11.65.82.82	28.8 5.8 8.78 8.02 122	24.82.82 25.12.82 18.12.82 18.13.82 18.13.82 18.13.83 18.	1.000 1.000	
15.54 29.56 79 0.273	y, inch	0. 0035 . 0085 . 0135 . 0185	0285 0335 0435 0485	.0535 .0785 .1035 .1285 .1585	2035 2535 3035 3535 4035	. 4535 . 5035 . 5535 . 6035 . 6535	7035 7535 8035	
6 6 6 6	1/n	0, 442 365 368 398 626 641	853 873 888 888 889	704 735 737 794	808 828 850 878 900	929 957 957 957 957	206 888 208 888 208 888 208 888	666
14. 09 29. 66 85 0. 249	y, inch	0. 0035 .0085 .0135 .0185	0285 0335 0435 0485	.0535 .0785 .1035 .1285	2035 2035 2535 3035 3535	. 4035 . 4535 . 5035 . 5535 . 6035	6535 7035 7535 8035 8535	. 9035
12, 55 t) 30, 00 91 0, 221	u/17	0, 431 .557 .595 .619 .635	25.9 15.5 15.5 15.5 15.5 15.5 15.5 15.5 1	######################################	8.22 8.22 8.49 8.74 8.97	920 940 957 983	996 998 998 1,000	
$P_{t\infty} \times 10^{-6}$ $p_{t\infty}$ (psia) T (° F) M	y, inch	0.0035 .0085 .0135 .0185	0285 0335 0335 0435 0485	.0535 .0785 .1035 .1285	. 1785 . 2035 . 2535 . 3035 . 3535	.4035 .4535 .5035 .5535 .6035	. 6535 . 7035 . 7535 . 8035 . 8535	

${\bf TABLE\ I.-MEASURED\ BOUNDARY-LAYER\ VFLOCITY\ PROFILES--Concluded}$

(d) x = 51.75 inches—Concluded

$P_x imes 10^{-6}$ P_{∞} (psia) P_{∞} (psia) M	32, 54 76, 25 89 0, 224	34. 78. 95 0.	78 02 238	36, 82 80, 29 104 0, 251		39, 32 81, 69 108 0, 263		41, 24 80, 84 94 0, 275	
, inch	u/V	y, inch	u/V	y, inch	u/V	y, inch	u/V	y, inch	u/V
), 0035	0.483	0, 0035	0. 489	0.0035	0, 495	0,0035	0, 500	. 0, 0035	0, 51;
, 0085	. 589	.0085	. 594	. 0085	. 597	.0085	, 602	. 0085	, 60
.0135	. 631	. 0135	. 633	.0135	. 636	0135	, 639	.0135	643
. 0185	. 655	. 0185	. 658	.0185	. 658	.0185	. 663	0185	. 66
. 0235	. 673	. 0235	. 675	. 0235	, 675	, 0235	. 680	. 0235	. 68-
0285	. 687	0285	. 689	, 0285	. 691	. 0285	693	. 0285	. 697
. 0335	. 698	. 0335	. 701	, 0335	. 702	. 0335	. 704	. 0335	. 709
. 0385	. 709	. 0385	. 711	. 0385	. 712	. 0385	. 713	. 0385	. 71
. 0435	. 717	. 0435	. 719	. 0435	. 722	. 0435	. 722	. 0435	. 72
. 0485	. 725	. 0485	. 726	. 0485	. 729	. 0485	. 730	. 0485	. 73
. 0535	. 732	. 0535	. 733	, 0535	. 735	. 0535	. 737	. 0535	. 74
, 0635	. 743	, 0635	. 745	. 0635	. 747	. 0635	. 748	. 0635	. 75
. 0735	. 754	. 0735	. 755	. 0735	. 756	. 0735	. 757	. 0735	. 760
. 0835	. 764	, 0835	. 764	. 0835	, 766	. 0835	. 767	. 0835	. 77
. 0935	. 772	. 0935	. 773	. 0935	. 775	. 0935	. 775	. 0935	. 77
. 1035	. 780	. 1035	. 781	. 1035	. 782	. 1035	. 783	. 1035	. 78
. 1285	. 798	. 1285	. 799	. 1285	. 800	. 1285	, 800	. 1285	, 80
1535	. 813	. 1535	. 815	. 1535	. 817	. 1535	. 818	, 1535	. 820
. 1785	. 828	. 1785	. 829	. 1785	, 832	. 1785	. 832	. 1785	. 833
2035	. 842	. 2035	. 843	. 2035	. 845	. 2035	. 847	. 2035	. 84
. 2535	. 869	. 2535	, 870	. 2535	. 870	. 2535	. 872	. 2535	. 87
. 3035 . 3535	. 892	. 3035	. 893	. 3035	. 894	. 3035	. 895	3035	, 89
. 4035	. 914 . 935	. 3535	. 915	. 3535	. 916	. 3535	. 916	. 3535	. 919
. 4535	. 935 . 951	, 4035 , 4535	, 935 , 952	. 4035	, 936	4035	. 937	. 4035	. 93
Total .		4000	. 982	. 4535	. 955	4535	. 954	. 4535	, 95
, 5035 , 5535	. 966 . 979	. 5035 . 5535	. 967	. 5035	, 968	. 5035	, 969	. 5035	. 97
6035	. 979	. 5535 6035	, 980	, 5535	. 981	. 5535	. 980	, 5535	. 98.
, 6535	. 989 . 994	. 6535	. 989	. 6035	. 989	6035	. 990	. 6035	. 99
. 7035	. 994	. 6555	. 994	6535	. 994	6535	. 994	. 6535	. 996
			. 996 	. 7035	. 997	. 7035	. 997	. 7035	. 998
. 7535	. 997	. 7535	. 997	. 7535	. 998			. 7535	. 998
. 8035	. 998	. 8035	. 997	8035	. 998			. 8035	. 996
. 8535	. 998	. 8535	. 998					. 8535	. 999
9035	, 998					į			

TABLE II.—MEASURED LOCAL SKIN-FRICTION COEFFICIENT

(a) x = 18.00 inches

	C/×10³	25 55 55 25 55 55 25 55 55	2.60	2.2.2.2 2.2.2.2 53.5.5.4 53.5.5.4		
67.26	P,×10−6	25.85.6 25.85.6 25.85.6 25.85.6	7.34	8 5 2 6 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8		
	W	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	. 153	188 108 158 158		
	C _f ×103	21226 2126 2126		22222 2886 2986 2986 2986	2. 61 2. 58	<u> </u>
48.81	P ₄ ×10→	8 E 8 E 8	3.79	55.4.5.5.5.5.5.5.5.5.5.5.5.5.5.5.5.5.5.	## 0 ## 12	
; ;	Ж	0. 107 153 188 198	601	152 187 218 218 107 152	. 187	
1	C;×103	22.50 27.50			999999 1817538	255255 25755
29.99	$R_t \times 10^{-6}$	3.3.3.4.5.0.3.0.0.0.0.0.0.0.0.0.0.0.0.0.0.0.0.0	5. 47	25 25 25 25 25 25 25 25 26 26 26 26 br>26 26 26 br>26 26 26 26 26 26 26 26 26 26 26 26 26 26 26 26 26 26 2	222 222 222 222 222 222 222 223	20.4.0.0.0.0.4 20.8.8.8.8.9
	W	0.153 187 219 246	569	931. 136. 189. 220. 220.	. 246 . 270 . 312 . 155	28.2 28.2 290 290 290 290
	C/×103	8.8.8.8.8 8.8.8.8.8 8.8.8.8.8.8 8.	5. 65	22.98 2.98 2.98	50 50 50 50 50 50 50 50 50 50 50 50 50 5	i
15.15	$P_x \times 10^{-6}$	2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,	12	2999999 88889 88889	999999 893999	
	N.	0. 225 . 247 . 270 . 270	314	220 246 269 291 313	247 247 269 292 314	
	C×103	20 20 20 20 20 20 20 20 20 20 20 20 20 2		20 # # # 20 # # # # # # # # # # # # # # # # # # #	25 25 25 25 25 25 25 25 25 25 25 25 25 2	
$p_{t\infty}$ (psia) 7.55	$P_x \times 10^{-6}$	1.248		1.20	1.1.2.2.1.2.3.2.2.1.3.3.2.2.1.3.3.3.2.2.2.3.3.3.3	i <u></u>
P _t	N	229 249 275 205	317	223 249 254 255 315	252 252 272 296 318	# !

85	C,X103	999999 88888 8888 8888
(psia) 75.28	P _x X10→	68 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8
$p_{t\infty}$	X	8.2.8.8.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2

TABLE II. -MEASURED LOCAL SKIN-FRICTION COEFFICIENT--Continued

				_		_		_	_					
	C/X103		2.60	2.50	2, 2, 2, 88 2, 88	9.55	. S. S.	25.5	2, 57	2, 57	93 0	2.0		
48.41	P _x ×10-6	6.37		 	6.28	\$ 000	16 18	12, 40	6.28	95. 95.	10 76	10.5	i	
	M	0 110	. 153	189	. 109	. 154	951	. 218	. 109	1.2	231	218		
	C/X103				5 5 5 6 12 6			2.46						2.65
28.55	P,×10-6	6.63	9.35	200 200 100 100 100 100 100 100 100 100	7.05 2.05	6.15	5.03	9.95	ا ا ا	8. 65	2 6	20.22	6, 15	5.05
	X	0.312	. 291	. 200	217	. 188	. 153	. 312	292	. 508	244	217	188	. 153
	C ₇ ×10 ³	2.93	2.83	9.7 0.10	25.2	2.84	2.83	2.80	2,76	2.70	2.85	2.83	2. 77	25.72
15.14	P. X 10 3	3.97	+ +	4. n 8. n	5.52	3.92	4.36	. 78 . 78	5.16	5. 40	68.8	4.35	1.81	5.15
	X	0. 223	. 247	277	315	. 220	. 247	. 271	. 293	. 312	.219	. 246	. 272	293
	C;X103				5 2	2.81	2.77	2.72	7.5	60.7				C ()
14.90	P _x ×10·4	3.90	4 8. !	7 Z	5.39	3.84	 96.	 [3]	0.01 3.01	9. 9/	3.86	4.27	4.	4, 4 8 8
:	M	0.218	. 245	6 5	. 313	.218	. 245	58.0	167.	010	.218	. 245	. 268	230
	$C_f \times 10^3$	3. 25	8. m:	- 35 15 mi	3.05	3.18	%; 6;	5. 5. 5. 5.		0.00		3.24	 	10 d m m
$p_{t_{\infty}}$ (psia) 7.68	P _x ×10→			38		1.99		2.46						2.5
p,	W	0. 223	247	96	.316	219	253	217	202.	.	. 220	.250	274	900

7	Pto (pera) or.49			70.11	
ж	P,x10-6	$C_f \times 10^3$	K	P.×10-6	C ₇ ×10 ³
0.218			0.110		
<u>38</u>			. 155		
Ŧ.	12, 56	2.39	38	16 98	
801.			. 110		
. 219		2.32	. 154	13.99	2, 43
28			<u>\$</u>		2.33
. 153	12.54	2. 37	. 110		
.107			. 155	14.05	2, 40
	_		. 189		

TABLE II.—MEASURED LOCAL SKIN-FRICTION COEFFICIENT.-Continued

	C/X103	02 22 22 22 22 23 24 24 24 24 24 24 24 24 24 24 24 24 24	2.52 2.52 4.52 4.52	22.22 24.45 24.45 24.45		C×103	22.22.22 22.22.23 21.22.23 21.23	228882 138882	25222 2017 2017 2017 2017 2017 2017	222222 222222 223222 24222 25222 25222 25222 25222 25222 25222 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 2522 252 2522 252	2 2 18 2 2 19 2 2 21 2 23	2.30 38.33
28.94	P2 × 10-6	7. 36 8. 97 10. 36 11. 54 12. 59	13 149 14 28 15 28 16 29 20 20 20 20 20 20 20 20 20 20 20 20 20	11. 48 12. 58 13. 50 14. 21	79.27	Pz×10-6	88.28.28 80.28.29 80.08.29 70.08.29	26.93 19.15 19.18 33.61	32.83.94 28.53.94 28.53.85 28.53.85	23. 55 13. 10 13. 34 32. 92	1888.88 4.48.88 8.88.88	13.64
	N	0. 156 . 190 . 220 . 247 . 271	292 314 154 190 220	247 272 283 818 818		M	0. 275 . 269 . 258 . 245 . 232	. 219 189 . 153 . 109 276	270 258 245 232 232	189 153 108 277 270	258 246 232 219 189	109
	C _f ×10³		88282 88282			C ₇ ×10³	#2828 20000	225.22 225.22 225.23	:	 	 	
15.01	P _x ×10-6	5.55 5.98 5.98 4.68 5.09 5.09 5.09 5.09 5.09 5.09 5.09 5.09	5.32 5.48 6.48 7.702 46	88.98 8.49 8.49 8.49 8.49 8.49	75.19	$P_x \times 10^{-6}$	13, 78 19, 30 19, 30 19, 20 20, 20	23. 47 13. 55 19. 25 23. 42				
	7	0.217 2.268 2.289 3.11	218 244 267 291 312	218 246 268 291 312	:	M	0.109 153 188 108 153	851 858 858 858 858 858 858 858 858 858				!
	C ₇ ×10 ³	999999 699999 8999	55 55 55 55 55 55 55 55 55 55 55 55 55	999999 9889999 9889999		C;X108	2. 32 2. 32 2. 27 2. 30	8285 6666				
15.10	$P_x \times 10^{-6}$	6.03 6.03 7.17 6.41 6.41 6.41	5. 41 6. 58 6. 58 7. 12 6. 58	6.69 6.69 6.61 6.61 6.61 6.61	67.59	P ₄ ×10 %	12, 22 17, 28 21, 04 12, 17 17, 36	21.09 17.38 17.38 1.18				!
	W	0. 244 218 269 293 313	218 244 268 291 313	. 221 . 245 . 269 . 292 . 314		<i>X</i>	0.108 1.154 1.108 1.154	188 108 153 189			 	
	C ₇ ×10³	8.88.89 9.28.89 9.28.89	668888 888886 888886	9899999		C/X103	2222 2337 238 238	228824 28824 28824	88 58 61 64	 	!	
11:1	P,×10-6	93.85 93.48 93.75 115	2. 82 3. 17 3. 46 3. 72 3. 92	26.88.88 27.25 27.25 27.25	48.25	$P_x \times 10^{-6}$	8.82 15.03 15.03 2.73 2.73	58.52 23.52 23.53	15.02	! !		
	X	0.225 .276 .317 .24S	221 250 273 297 314	222 250 273 296 316		X	0. 109 . 153 . 188 . 218	153 108 108 153	2188	·		
	C:X163	8.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2	22.23	2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2		C _f ×10³	256844 256844	998888 8888	\$\$\$\$\$\$ \$35533	22222 442333 4423333	17.7	
$p_{i,\infty}$ (psia) 7.76	P,×10-6	9.82.82 9.82.82 9.82.83 9.82.83 9.82.83	25 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8 8	2.5.5.5.5. 2.5.2.8. 2.5.2.8. 2.5.2.9.	(psia) 29.29	P _x ×10-6	14. 49 13. 59 12. 64 10. 31	6 2 4 4 4 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	11. 54 10. 91. 00 14. 31 14. 52	23.11.12.62 10.246 28.33.246 38.33.246	7.25	
$p_{i\alpha}$	· .	0. 225 249 272 300 316	215 254 274 318	223 252 276 297 315	$p_{i_{\varpi}}$	<i>X</i>	0. 313 292 269 245 245 218	188 313 292 292	245 219 188 153 314	292 269 244 218 188	. 152	

TABLE II. MEASURED LOCAL SKIN-FRICTION COEFFICIENT—Concluded

79,04	P,X10-4 C;X103	2000	39.18 41.10 42.83 43.25 17.80 2.38	25.11 2.26 30.51 2.20 35.14 2.15 37.04 2.13 38.98 2.13	40.83 42.31 42.98 17.69 22.45 25.01	
	W.	0. 111 156 190 223 233	251 251 273 111 111	221 221 234 234 247	272 277 111 136	. 191 . 220 . 233 . 247
	C ₇ X10 ³	22.24 22.24 1.17	22222 5222 5222 722	91 64		
74.68	R _z ×10-6	25. 28 26. 28 26. 28 26. 28 26. 22	25.25 25.28 25.38 25.38 25.38	30.46		
	W	0.109 154 189 109	81. 82. 82. 83. 83. 83. 83.	. 189		
	C/X108	999999 582248	2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,2,			
67.85	P.×10-6	22, 67 27, 45 16, 05	27. 45 15.93 22. 48 27. 35		; ! !	
	W	0.109 .155 .189 .108	. 189 . 109 . 154 . 189		:	
	C ₇ ×10 ³	문중 있 었 건강점점건	444444 4888 8888 8888	22.3	!	
48.03	R _x ×10-5	1.22 1.22 1.23 1.25 1.25 1.25 1.25 1.25 1.25 1.25 1.25	15, 95 19, 48 22, 32 11, 22 15, 93	19. 46 22. 32		
	M	0. 109 - 154 - 189 - 219	189 189 218 201 108	. 189		
	C/X10 ³	조건성성성 조건성성성	स्ट्रह्म इंट्रह्म	ଅଧ୍ୟୟୟ ଅଧିକ୍ରୟ		
$p_{i,\infty}$ (psia) 7.61	P _x ×10-5	25.44.4.5.5.5.5.5.5.5.5.5.5.5.5.5.5.5.5.	3.71 4.12 5.14 5.16 5.16	3 70 4 19 4 53 5 19		
	M	0, 226 . 254 . 277 . 300	224 251 274 300 319	255 277 297 320		

	ž.	

TABLE III.—SUMMARY OF MAJOR PARAMETERS OBTAINED FROM MEASURED BOUNDARY-LAYER PROFILES

 $p_{t\infty}$ Mδ* r. in. $P_x \times 10^{-6} \theta_1$ in. δ^* , in. δ , in. C_{i} θ psia 0, 308 7, 94 , 215 15, 17 , 241 14, 81 , 264 14, 83 , 288 14, 84 15, 75 15, 75 15, 75 15, 75 15, 75 15, 75 1, 43 1, 97 2, 15 2, 33 2, 51 0. 0331 . 0305 . 0301 1, 378 1, 390 1, 372 1, 378 0, 0456 | 0, 296 0.00346. 263 . 257 . 247 . 245 . 00324 . 00318 . 00313 . 0424 .0292 . 00309 03971.360.307 14.86 .152 28.41 .187 28.65 .216 28.94 .243 29.25 15, 75 15, 75 15, 75 15, 75 15, 75 2, 64 2, 65 3, 24 3, 72 4, 16 . 243 . 242 . 240 . 0288 . 0283 . 0276 0395 00305 . 0392 . 0387 $\frac{1.385}{1.402}$. 00306 . 00294 . 0265 . 0262 233 228 1. 381 1. 374 00286 0360 00282 29, 37 30, 73 30, 52 78, 45 78, 83 15, 75 15, 75 15, 75 15, 75 15, 75 . 265 . 289 . 308 . 107 4 56 0258 0348 1 365 .00278 . 0251 . 0245 . 0247 . 0344 . 0328 . 0340 1, 365 1, 371 1, 339 1, 377 . 00273 . 00271 . 00271 . 00273 5, 13 5, 48 5, 18 $\frac{219}{207}$.00253. 188 8.82.022603011.33227, 75 27, 75 27, 75 27, 75 27, 75 27, 75 7, 77 8, 23 7, 93 2, 52 2, 53 2, 53 . 0535 . 0733 . 453 1, 379 . 00309 . 309 . 307 . 307 . 214 . 216 . 447 . 447 . 410 . 415 1, 381 1, 410 .05290729 0722.00309.0518.00309 14. 75 14. 87 3. 33 3. 26 . 0480 1, 365 1, 377 . 00293 . 00294 27, 75 27, 75 27, 75 27, 75 27, 75 27, 75 15, 21 15, 01 15, 07 15, 01 14, 94 3, 47 3, 83 4, 17 4, 50 4, 70 0475 0651 . 402 . 403 . 398 . 394 1.371 00290 . 0476 . 0467 . 0457 1, 380 1, 365 1, 359 .00285 .00282 .00279 0650 0634 $288 \\
310$ 06210617 . 393 1, 362 .0027727, 75 27, 75 27, 75 27, 75 27, 75 27, 75 14, 94 14, 95 29, 51 29, 45 29, 46 . 00277 . 00277 . 00275 4, 62 4, 67 . 309 . 308 . 152 . 186 . 217 1, 361 1, 361 1, 336 .0454 .0618. 391 4, 90 5, 82 6, 74 . 384 . 377 . 371 0438 0596 . 0431 . 0416 . 0576 . 0561 1, 336 1, 349 . 00269 . 00263 27, 75 27, 75 27, 75 27, 75 27, 75 27, 75 . 243 29, 10 . 242 29, 97 . 242 29, 45 . 266 29, 62 .00259 .0408 0546 370 1, 338 . 371 . 366 . 364 0413 1, 358 1, 336 . 00259 . 00260 8, 09 9, 00 .00256.0408 0544 1, 333 30, 32. 0402 0531. 363 1. 321 .00252.310 | 30, 63 .310 | 30, 83 .309 | 30, 11 .108 | 85, 28 .152 | 84, 93 27, 75 27, 75 27, 75 27, 75 27, 75 27, 75 9, 60 . 0395 0525 . 00250 . 359 9, 49 9, 21 10, 04 . 361 . 360 . 367 . 346 , 00251 , 00252 , 00249 , 00239 0392.05291, 349 1, 337 . 0532 . 0520 . 0485 0398 . 0390 . 0370 1. 333 14.0839, 75 39, 75 39, 75 39, 75 39, 75 . 217 . 243 . 267 . 288 . 310 14, 79 14, 79 14, 87 14, 89 14, 90 4, 91 5, 55 5, 96 . 0650 . 0630 . 0618 0868 556 1, 335 00269 .0840 . 549 . 539 1. 333 1. 332 . 00266 . 00262 0607 . 542 . 527 6 36 0811 1, 336 .00260.0604 1. 326 00258 . 151 | 28, 72 . 187 | 29, 03 . 217 | 29, 26 . 244 | 29, 75 . 290 | 30, 35 6, 74 8, 31 9, 56 10, 68 12, 53 1, 326 1, 310 . 00269 . 0596 . 0790 516 0.583 .0764.00252. 0566 . 0561 . 0547 . 0753 . 0729 . 0715 1. 330 .00248 .00245 .00240 . 503 1. 300 1. 307 39, 75 39, 75 39, 75 39, 75 39, 75 39, 75 39, 75 311 107 152 187 221 30, 55 73, 39 74, 57 75, 87 78, 80 13, 45 12, 25 17, 34 21, 29 25, 82 1, 301 1, 316 1, 293 0549 0705 498 00238 . 488 . 483 . 473 , 00241 , 00241 , 00231 , 00226 . 0538 . 0516 . 0708 . 0667 0495 0638 1.289 0497 282 . 00221 39, 75 39, 75 39, 75 39, 75 51, 75 $\begin{array}{c} 235 \\ 247 \\ 260 \\ 272 \\ 317 \end{array}$ 78, 34 78, 22 79, 82 80, 51 . 00221 26, 21 . 0497 . 0639 . 483 28, 11 28, 86 30, 29 . 0625 . 0616 . 0615 . 476 . 470 . 469 1, 270 1, 275 1, 260 1, 353 . 00219 . 00218 . 00217 0492

.0483 .0488

.0887

.1200

. 743

4, 78

TABLE III.—SUMMARY OF MAJOR PARAMETERS OBTAINED FROM MEASURED BOUNDARY-LAYER PEOFILES -- Continued

r, in.	М	$p_{t\infty}$ psia	P _x ×10−6	θ, in.	δ*, in.	δ, in.	$\frac{\delta^*}{\theta}$	Cf
51, 75	0. 221	14, 84	6, 48	0.0802	0. 1067	0, 693	1. 330	
51.75	. 248	14, 89	7, 16	. 0781	. 1047	. 689	1.341	
51.75	. 274	14, 90	7. 92	. 0774	. 1018	. 676	1.319	
51.75	295	15, 13	8, 52	. 0751	. 1001	. 665	1. 333	
51, 75	. 317	15, 20	9.05	. 0751	, 1000	. 669	1. 332	
51. 75	. 155	29, 66	8, 96	0734	0972	, 655	1. 324	
51.75	. 190	29. 81	10, 95	. 0723	0954	. 655	1.320	0.00255
51, 75	. 221	30, 00	12, 55	0704	0925	. 643	1. 314	00250
51.75	. 249	29, 66	14, 09	0698	. 0909	, 636	1.302	. 00246
51, 75	. 273	29, 56	15, 53	0689	. 0893	. 643	1, 296	. 00242
51. 75	. 295	29, 76	16, 61	. 0677	. 0884	. 629	1.306	. 00239
51, 75	. 319	28, 47	16, 94	0688	0885	. 633	1. 286	. 00238
51, 75	. 110	84, 01	18, 89	. 0676	. 0871	. 633	1, 288	. 00234
51, 75	. 157	85, 15	26, 75	. 0630	. 0808	. 611	1. 283	. 00222
51, 75	. 193	√0, 49	30, 55	. 0622	0791	. 595	1.272	. 00219
51, 75	. 224	6, 25	32, 54	. 0633	. 0807	, 609	1, 275	. 00217
51, 75	. 238	78, 02	34, 78	. 0634	0802	. 609	1, 267	. 00215
51, 75	. 251	s0, 29	36, 82	. 0629	. 0791	. 603	1.258	. 00214
51, 75	. 263	51, 69	39, 32	. 0625	0792	. 603	1, 267	. 00212
51, 75	. 275	50, 84	41, 24	. 0606	. 0771	. 599	1.272	. 00210

TABLE IV. COMPUTED FRICTIONAL RESIST-ANCE LAW

C_f	P_{x}	Cy	P. e	$\frac{\delta^*}{\theta}$
·				U
0.0015	7.314×105	0, 00169	616450,	1, 216
, 0016	3, 918	. 00183	359220.	1,224
.0017	2.279	. 00193	219360.	1, 233
.0018	1. 366	. 00205	139670.	1, 241
. 0019	8, 513×107	. 00216	92100,	1, 249
. 0020	5, 487	. 00229	62700.	1, 257
. 0021	3, 636	. 00241	43852.	1, 266
. 0022	2.481	. 00253	31412.	1, 274
. 0023	1. 734	. 00266	23023.	1, 281
. 0024	1. 239	. 00278	17196.	1, 289
. 0025	8, 995×106	. 00290	13053.	1, 297
. 0026	6, 662	. 00303	10086.	1, 305
. 0027	5, 020	. 00315	7901.	1.312
. 0028	3.826	. 00327	6265,	1.320
. 0029	2.955	.00340	5028.	1, 327
. 0030	2.314	. 00353	4053.	1, 334
. 0031	1.830	. 00365	3343.	1.342
. 0032	L. 463	. 00378	2766.	1, 350
. 0033	1, 182	. 00390	2308.	1, 357
		İ		1.364
. 0034	9.621×10^{5}	. 00403	1941.	1, 371
. 0035	7.913	. 00416	1645.	
. 0036	6, 535	. 00429	1401.	1, 378
. 0037	5, 444	. 00442	1203.	1.386
. 0038	4, 561	. 00455	1038.	1, 393
. 0039	3. 853	. 00468	901.4	1, 401
. 0040	3,273	. 00481	787, 2	1, 407
.0041	2.794	. 00494	689, 5	1, 415
. 0042	2.401	, 00506	607, 4	1, 422
. 0043	2.069	. 00520	537, 5	1, 429
. 0044	1.795	. 00532	177.3	1, 436
. 0045	1, 560	. 00546	426.0	1, 443

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